

Coherent X-ray Diffractive imaging, Applications and limitations



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T. Beetz

- Coherent Diffractive imaging
- Phase-retrieval problem
 - Phase retrieval by projections
 - Examples
- Experiment at the Advanced Light Source
- Radiation Damage Limit
- Atomic-resolution imaging biological macromolecules
- Comparison of techniques

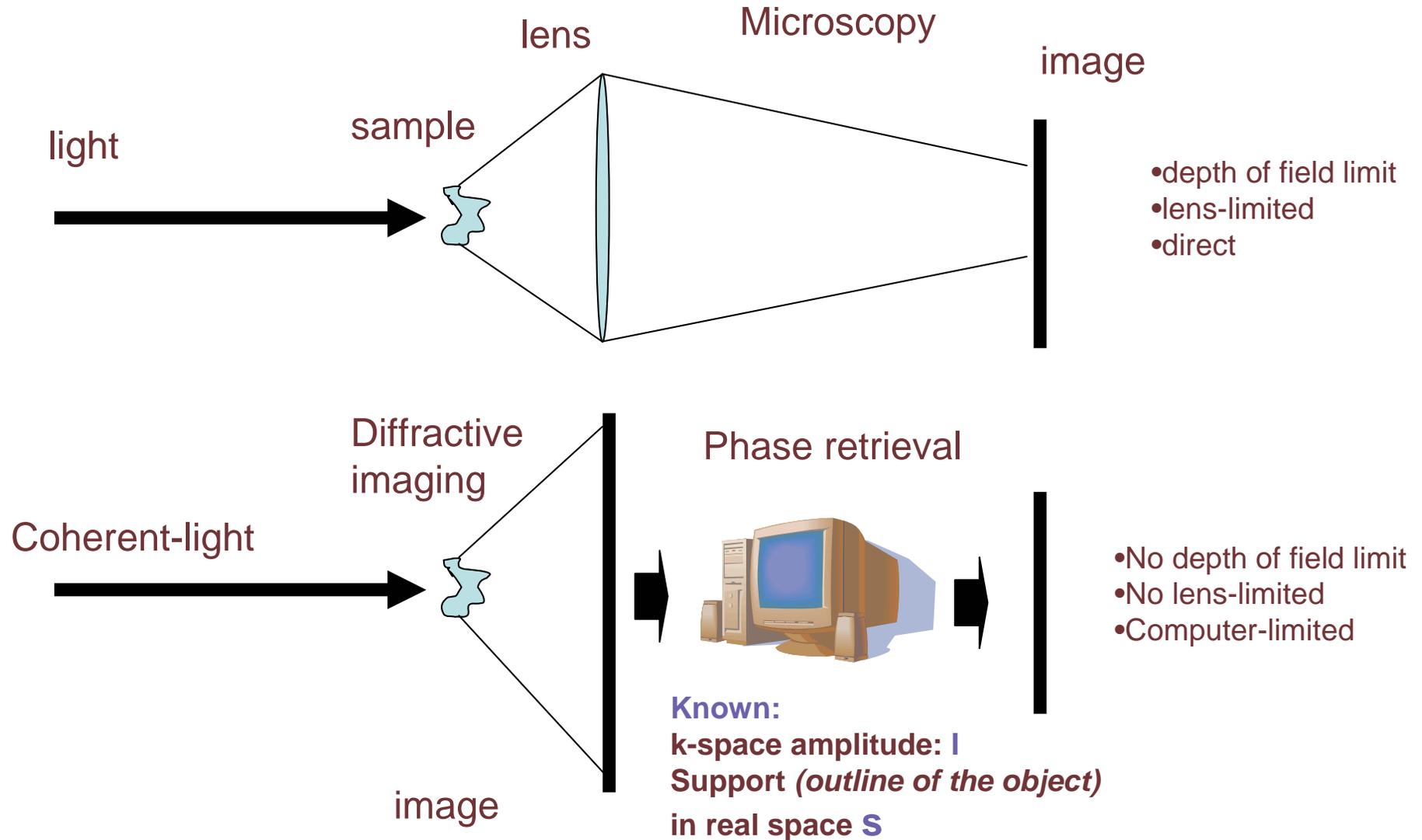


This work was performed under the auspices of the U.S. Department of Energy by the Lawrence Livermore National Laboratory under Contract No. W-7405-ENG-48 and the Director, Office of Energy Research, Office of Basics Energy Sciences, Materials Sciences Division of the U. S. Department of Energy, under Contract No. DE-AC03-76SF00098. SM acknowledges funding from the National Science Foundation. The Center for Biophotonics, an NSF Science and Technology Center, is managed by the University of California, Davis, under Cooperative Agreement No. PHY0120999.

UCRL-PRES-201069

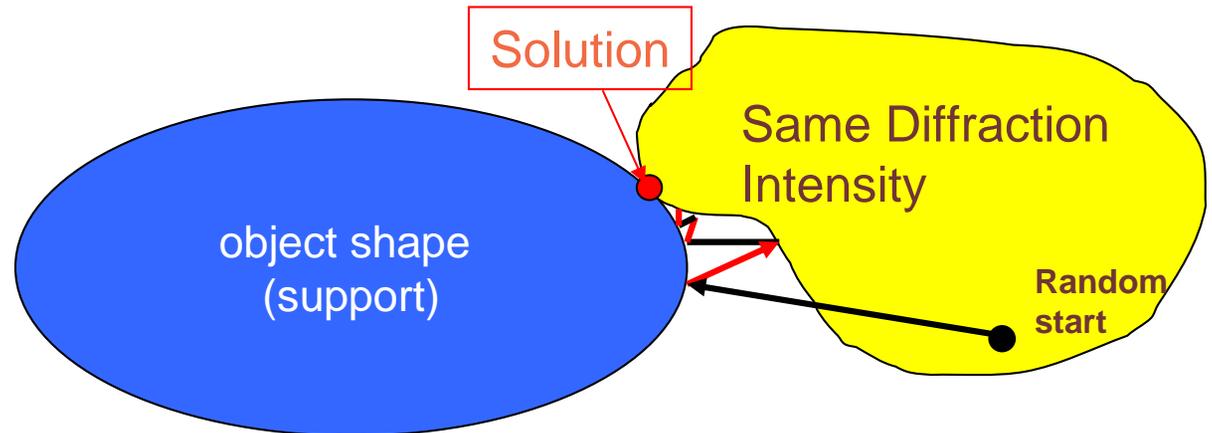
Coherent Diffractive imaging

phase-retrieval problem



Phase retrieval by projections

By projecting back and forth between feasibility sets, we can find the common solution (intersection of two sets)



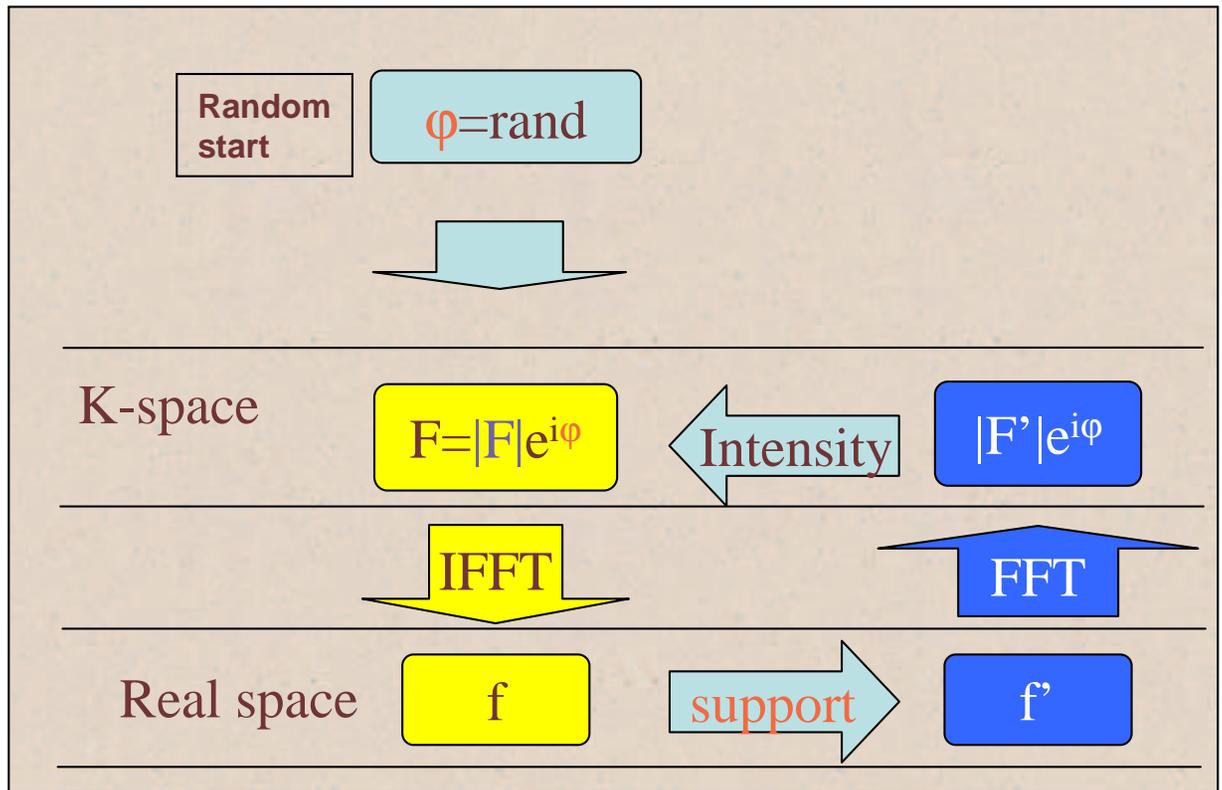
Known:

- k-space amplitude: $|F|$
- Support s in real space: $f=0$ for $s=0$

Unknown:

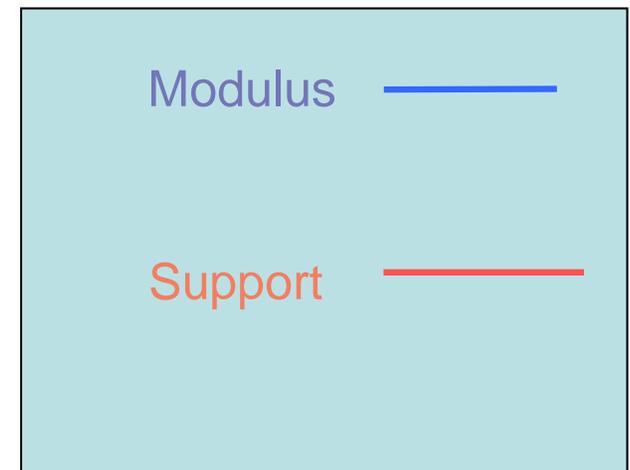
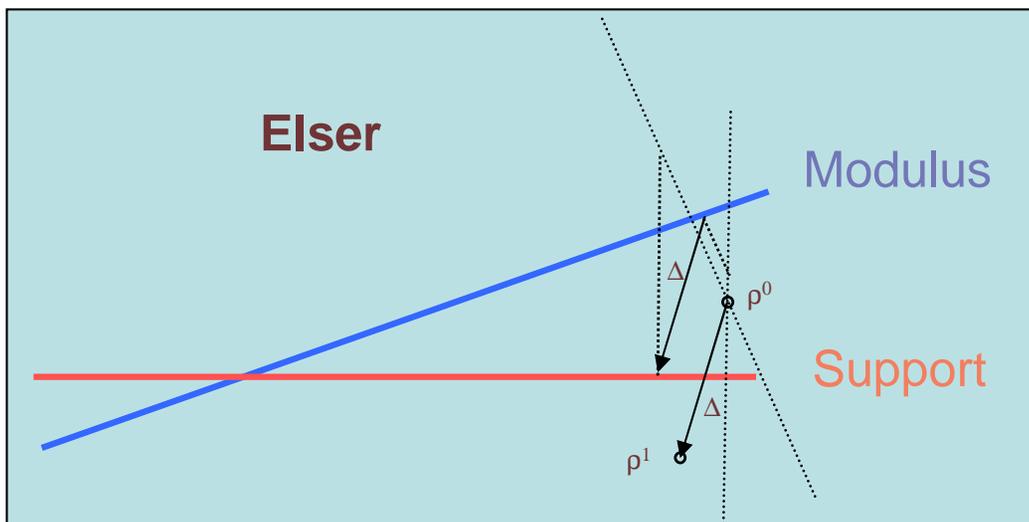
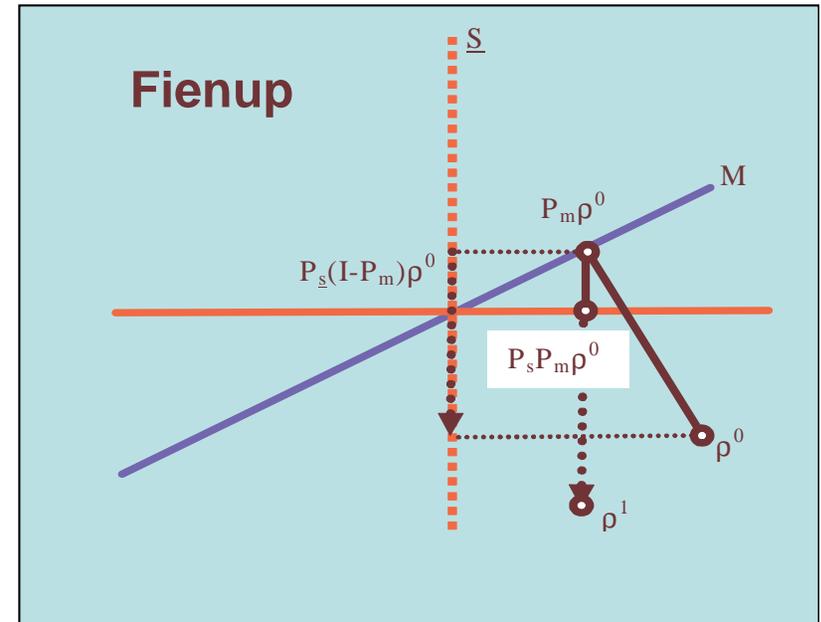
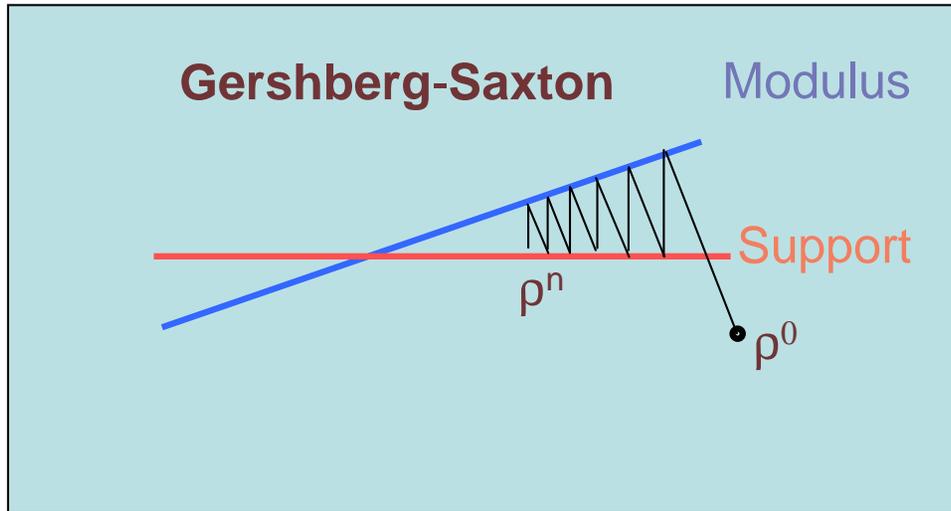
- Phase in k-space ϕ
- Image inside the support

$f=?$ For $s=1$

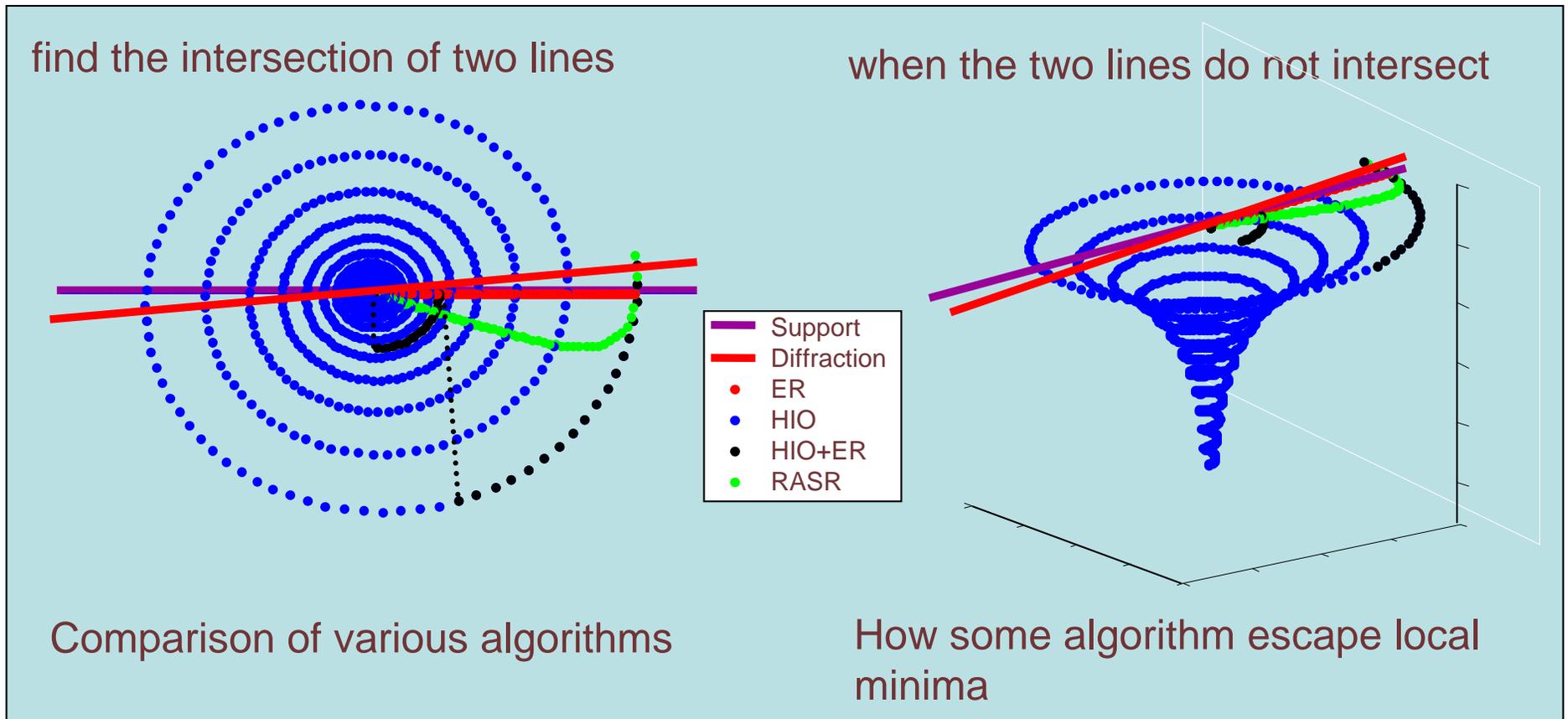


Projection algorithms

By projecting back and forth between feasibility sets, we can find the common solution (intersection of two sets)



Example:



- (ER) R. Gerchberg, W. Saxton, Optik 35, 237 (1972).
- (HIO) J. R. Fienup, J.R. (1982) . Appl. Optics, 21, 2758-2769
- (Diff. Map) V. Elser, J. Opt. Soc. Am. A 20, 40-55 (2003)
- (ASR,HPR, RASR) H. Bautschke, P.L. Combettes & D. R. Luke, (2002) . J. Opt. Soc. Am., 19, 1344-1345

Many error metrics

- real space error (stuff outside support)
- reciprocal space error, (R_{fact})
- difference between reconstructions

Reciprocal space error

$$R^n = \frac{\sum_{\mathbf{k}} \left| |\tilde{\rho}_{\text{obs}}(\mathbf{k})| - |\tilde{\rho}^n(\mathbf{k})| \right|^2}{\sum_{\mathbf{k}} |\tilde{\rho}_{\text{obs}}(\mathbf{k})|^2}$$

Real space error

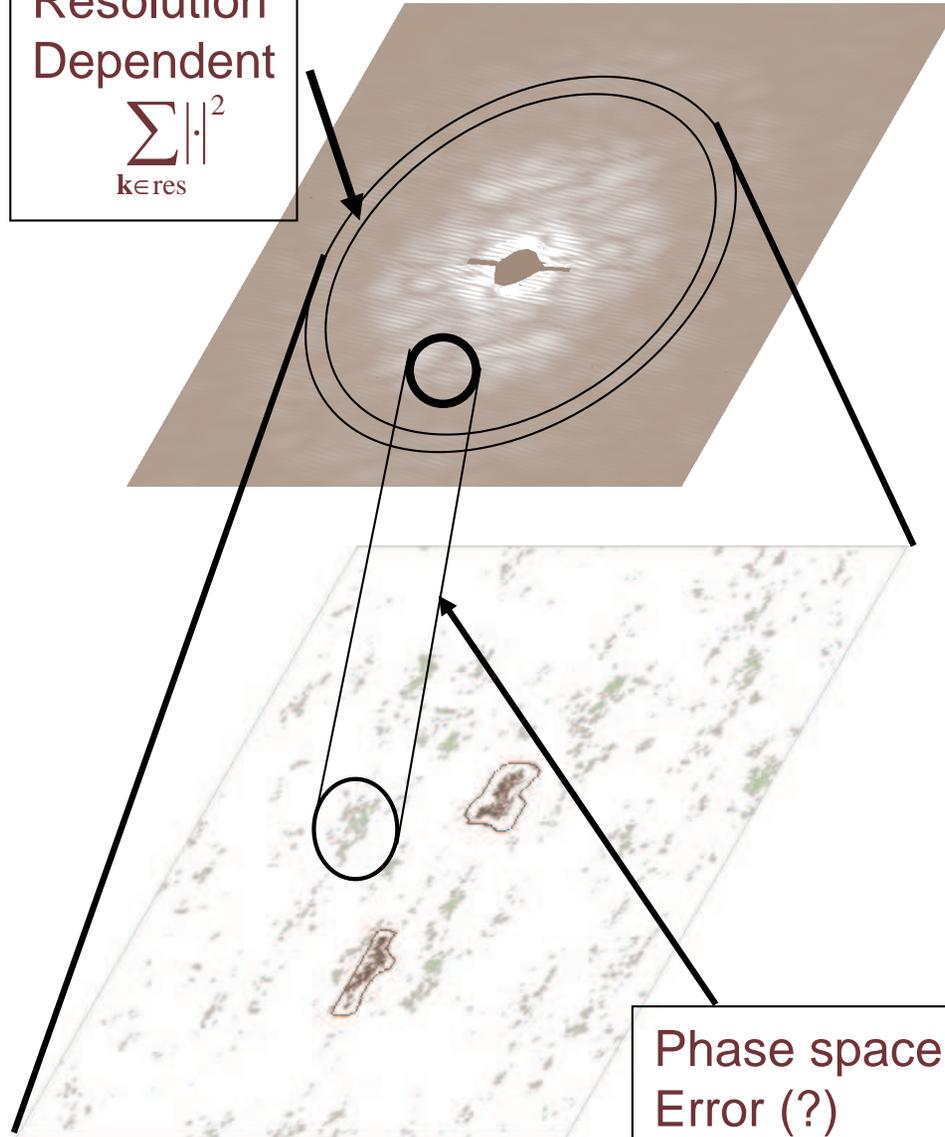
$$\epsilon_s^n = \frac{\sum_{\mathbf{x}} \left| \mathbf{P}_s \rho^n(\mathbf{x}) - \rho^n(\mathbf{x}) \right|^2}{\sum_{\mathbf{x}} \left| \mathbf{P}_s \rho^n(\mathbf{x}) \right|^2}$$

Error between reconstructions

$$\epsilon^{1,2} = \frac{\left\| \rho^1 - \rho^2 \right\|^2}{\left\| \frac{1}{2} (\rho^1 + \rho^2) \right\|^2}$$

Resolution
Dependent

$$\sum_{\mathbf{k} \in \text{res}} |\cdot|^2$$



Parseval

$$\begin{aligned} \|\cdot\|^2 &= \sum_{\mathbf{x}} |\cdot|^2 \\ &= \frac{1}{N} \sum_{\mathbf{k}} |\tilde{\cdot}|^2 \end{aligned}$$

In the
language of
projectors
and sets

Constraints errors

$$\epsilon_m^n = \frac{\left\| \mathbf{P}_m \rho^n - \rho^n \right\|^2}{\left\| \mathbf{P}_m \rho^n \right\|^2}$$

$$\epsilon_s^n = \frac{\left\| \mathbf{P}_s \rho^n - \rho^n \right\|^2}{\left\| \mathbf{P}_s \rho^n \right\|^2}$$

Phase problem not unique!

One family of homometric structures (Hoseman and Bagchi, Acta Cryst 7, p. 237 (1954)) may be generated using the result that.....

$$\rho_1(r) = l(r) * m(r)$$

and

$$\rho_2(r) = l(r) * m^*(-r)$$

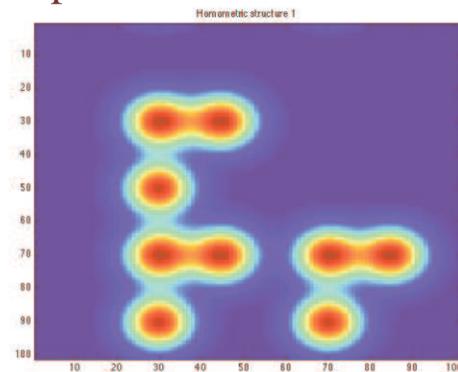
have same Fourier modulus $|R(u)|$, since

$$R_1(u) = L(u)M(u) \text{ and } R_2(u) = L(u)M^*(u)$$

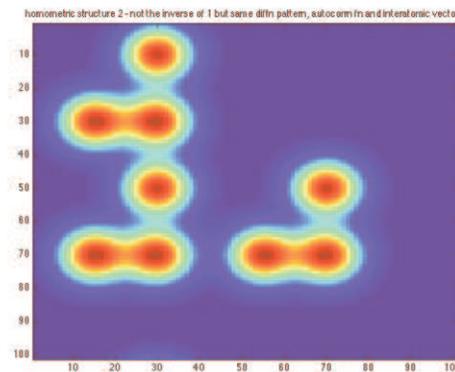
Homometric structures.

If $l(r)$ is a lattice and $m(r)$ a molecule, then m, m^* are enants, but ρ_1, ρ_2 are not enants.

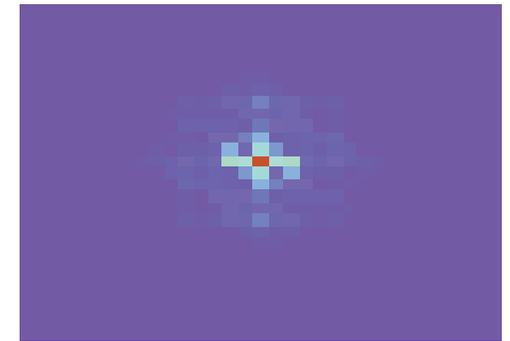
Example:



Homo1



Homo2



Fourier Mod of either

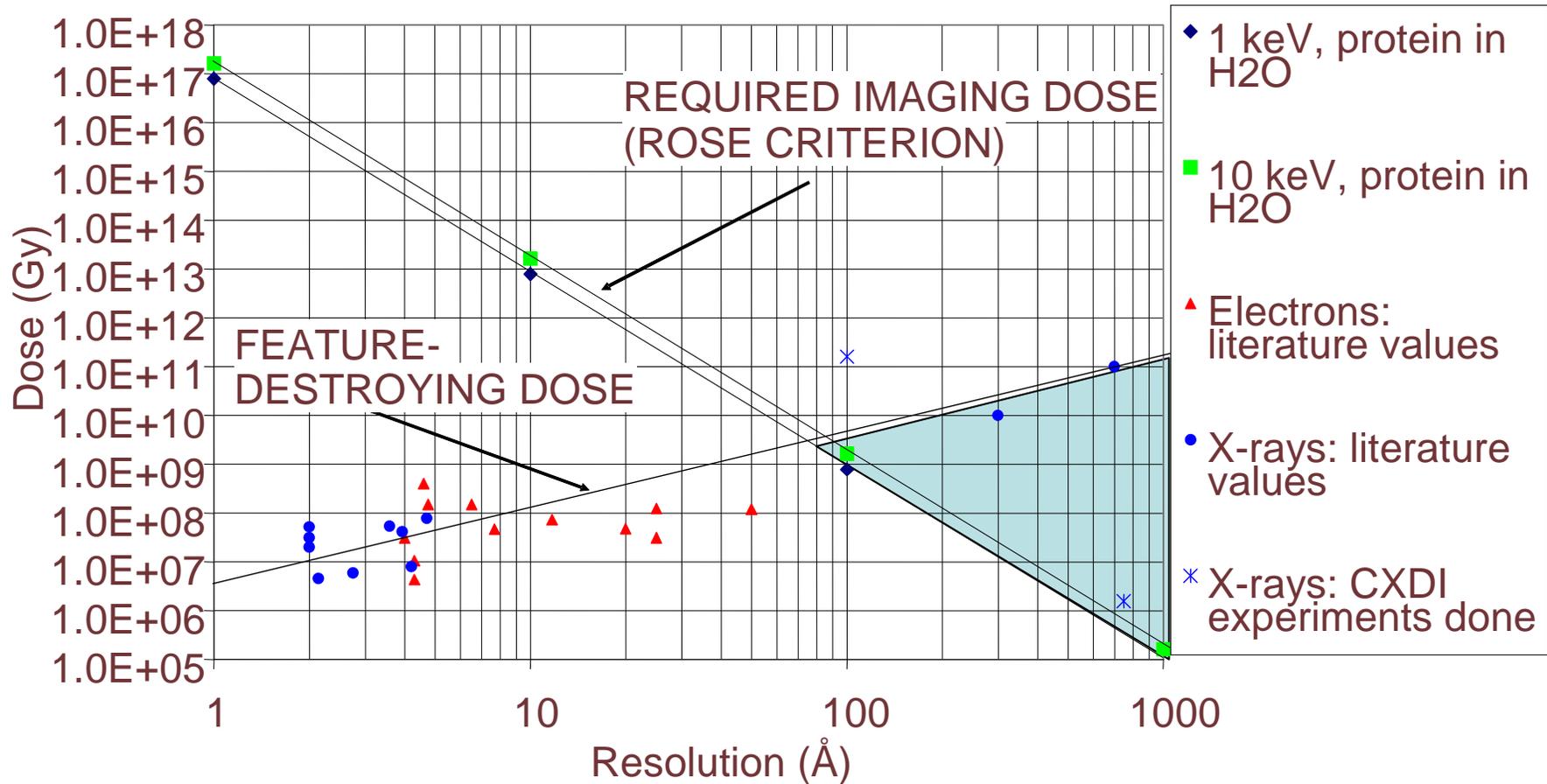
Note: Homo1 is not the inverse (enantiomorph) of Homo2.

Conclusion: HiO could not distinguish these unless tight support provided.

Distinguished by
Multiple scattering
ELNES \square

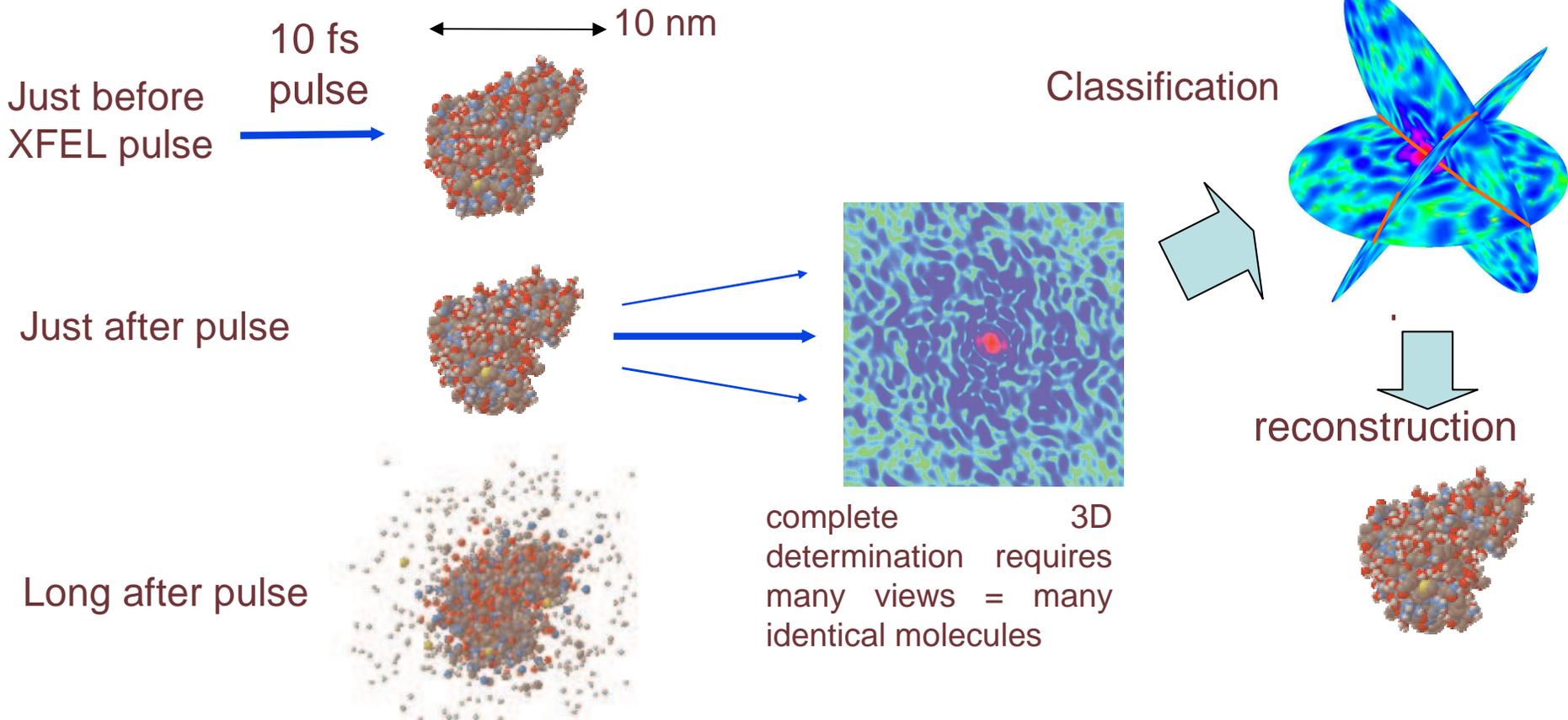
RADIATION DAMAGE LIMITED DOSE-RESOLUTION CURVES

M. R. Howells



Atomic-resolution imaging of virtually any biological macromolecule will become possible with XFELs

Diffraction from a single molecule:

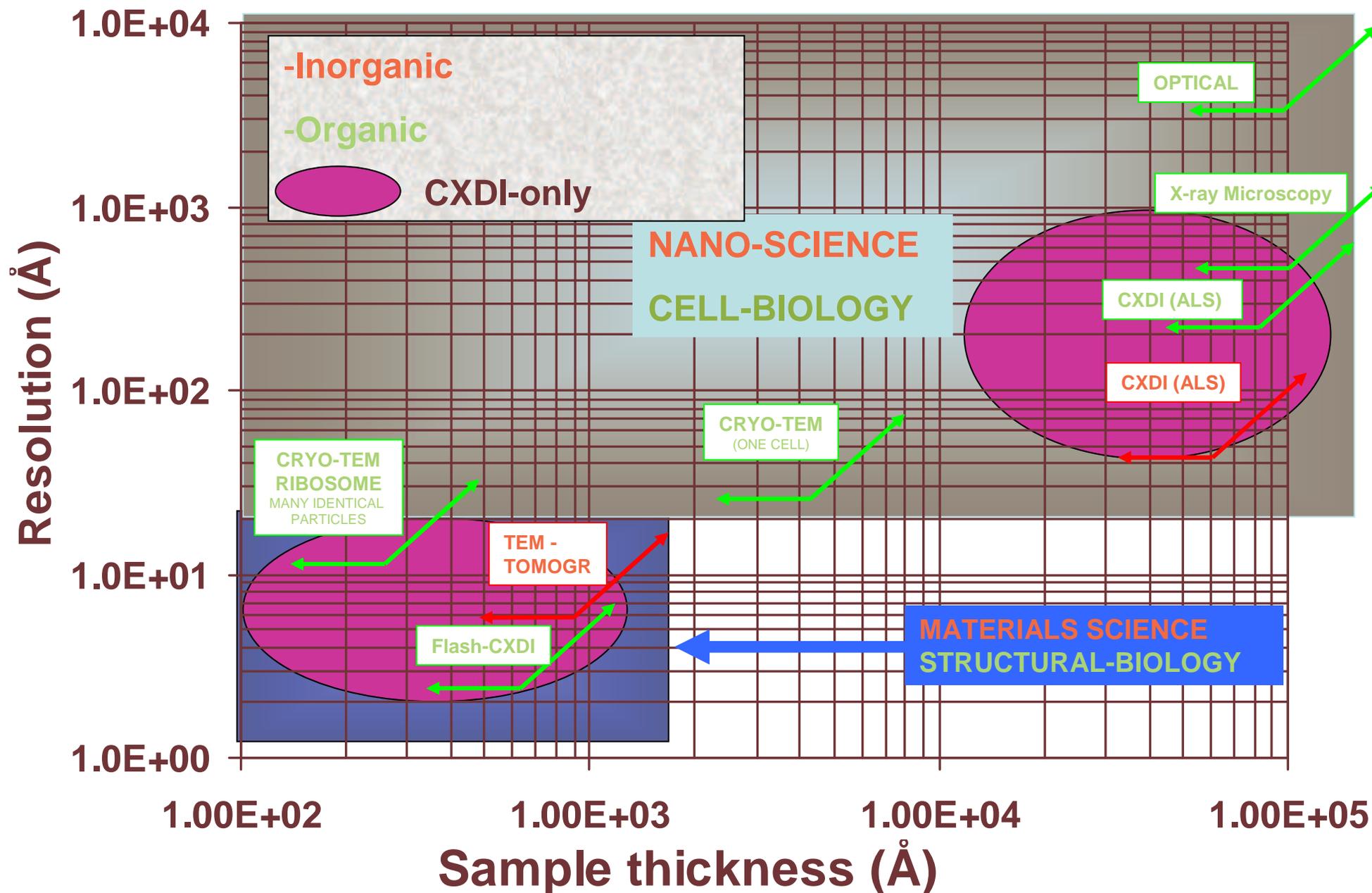


Many critical issues need to be solved:

- details of short-pulse photon-matter interactions
- image reconstruction
- sample manipulation and injection

Planned damage experiments at SPPS (Stanford), and TTF (Hamburg)

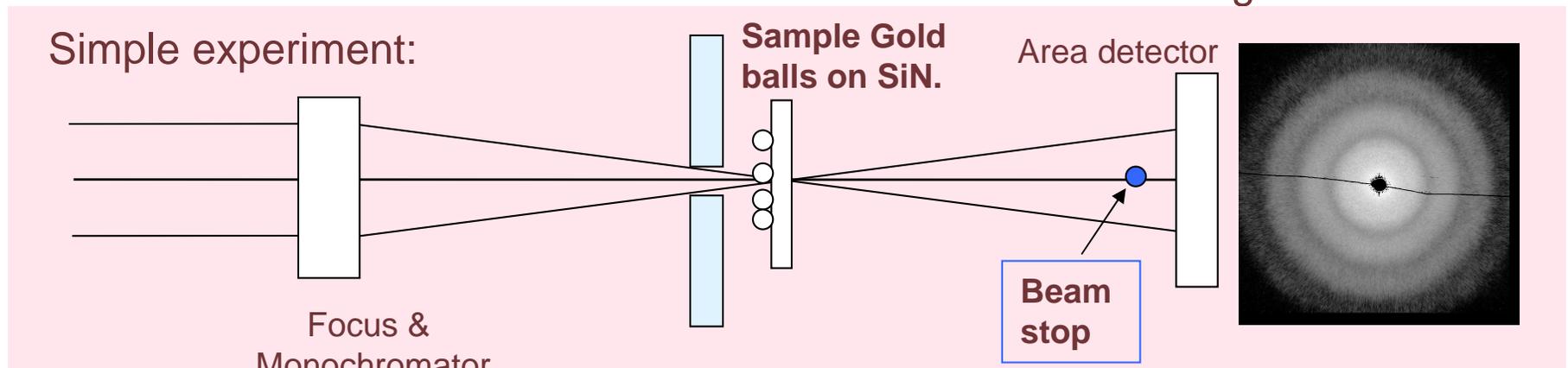
Comparison of techniques



Experiment

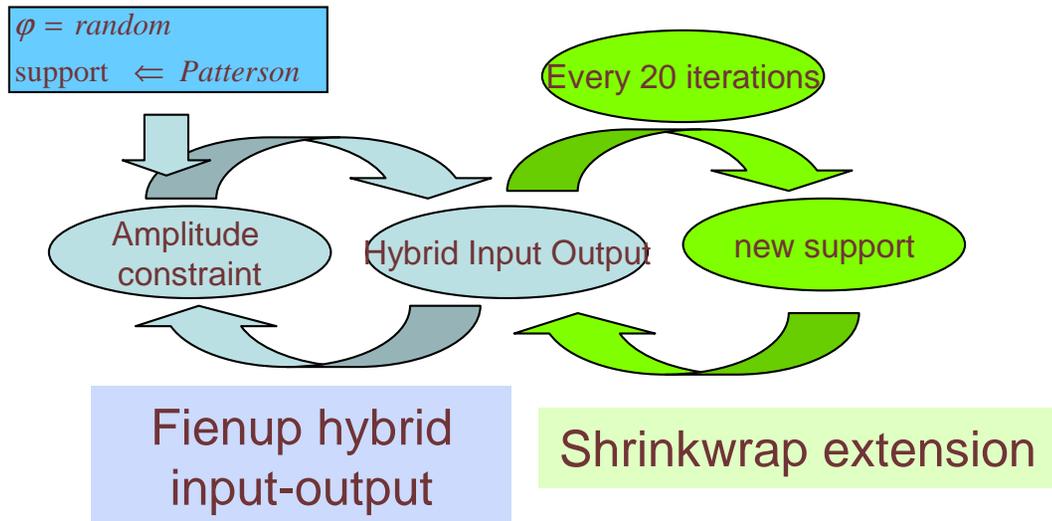
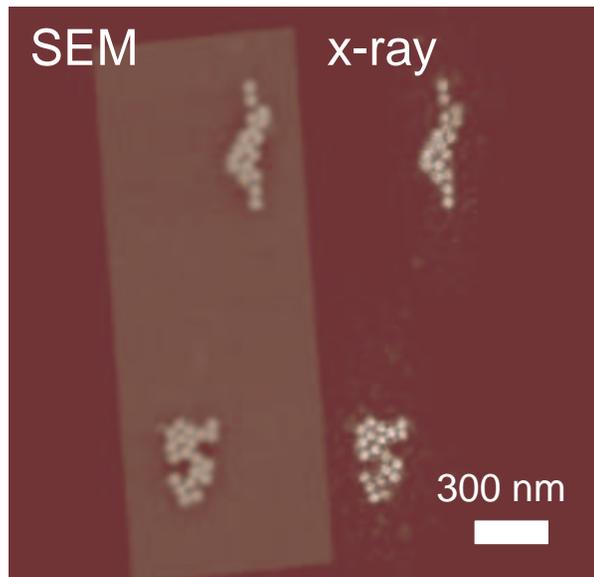
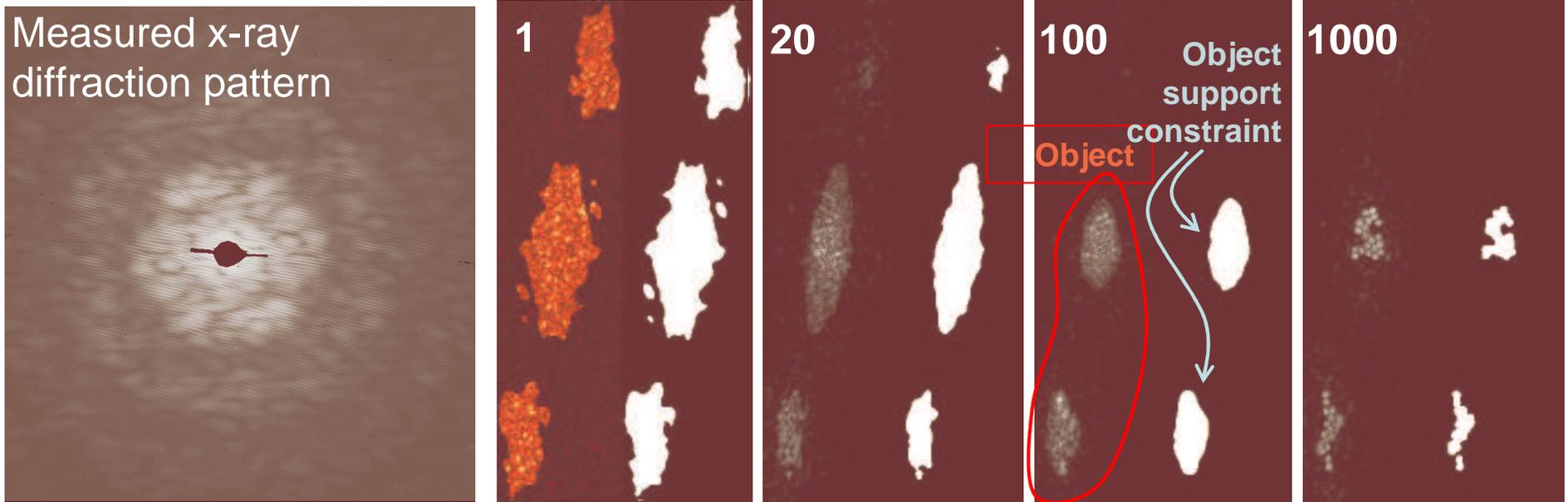
Layout of the diffraction chamber
at BL 9.0.1 at Advanced Light Source, LBL

Particle size of 50 nm scatters to 7 degrees at $\lambda = 2$ nm

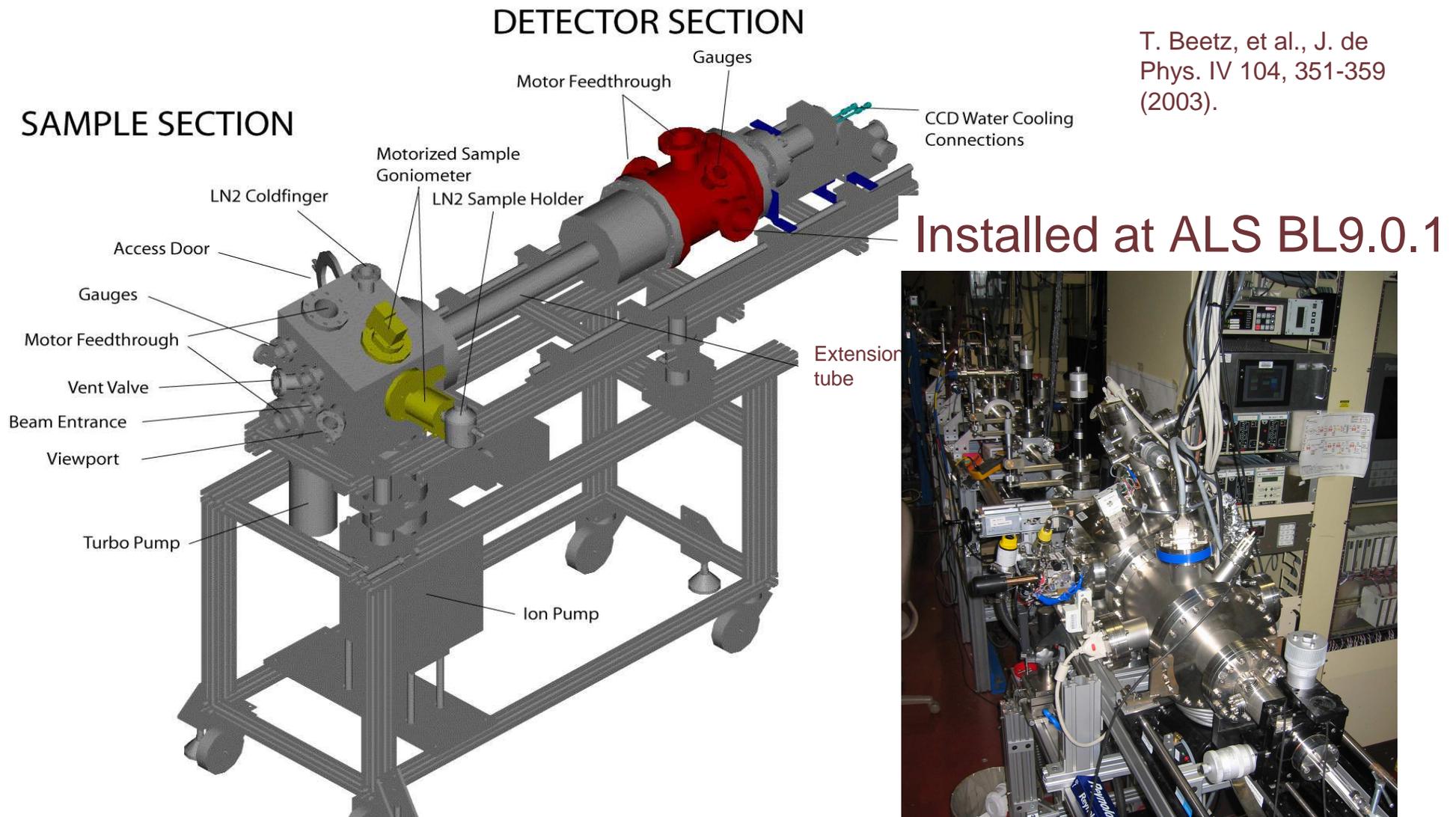


[Proc. SPIE 4783](#), 65 (2002)

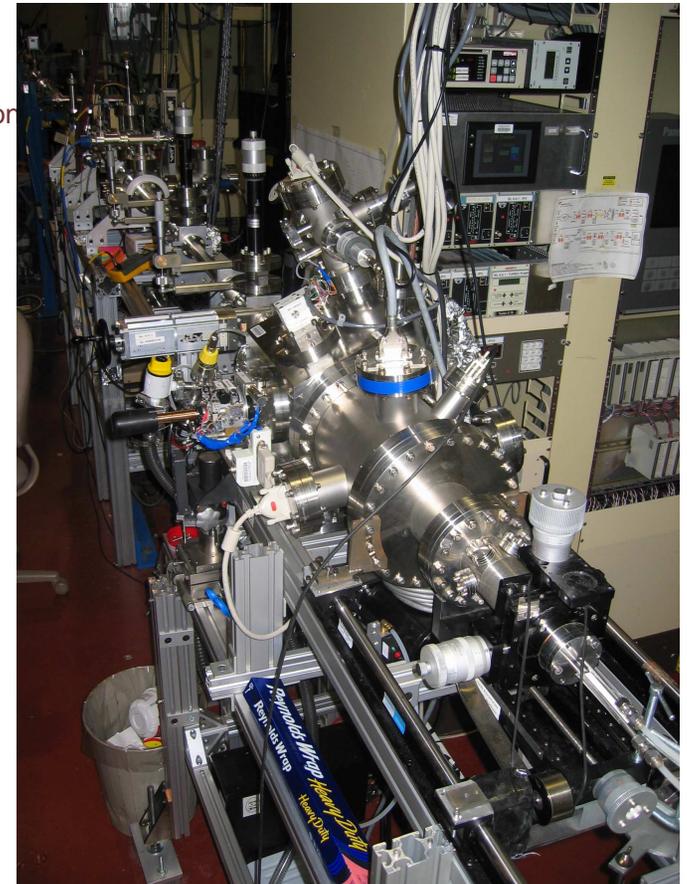
We have eliminated the need of the known shape of the object



The Stony Brook diffraction chamber allows accurate sample rotation and data acquisition

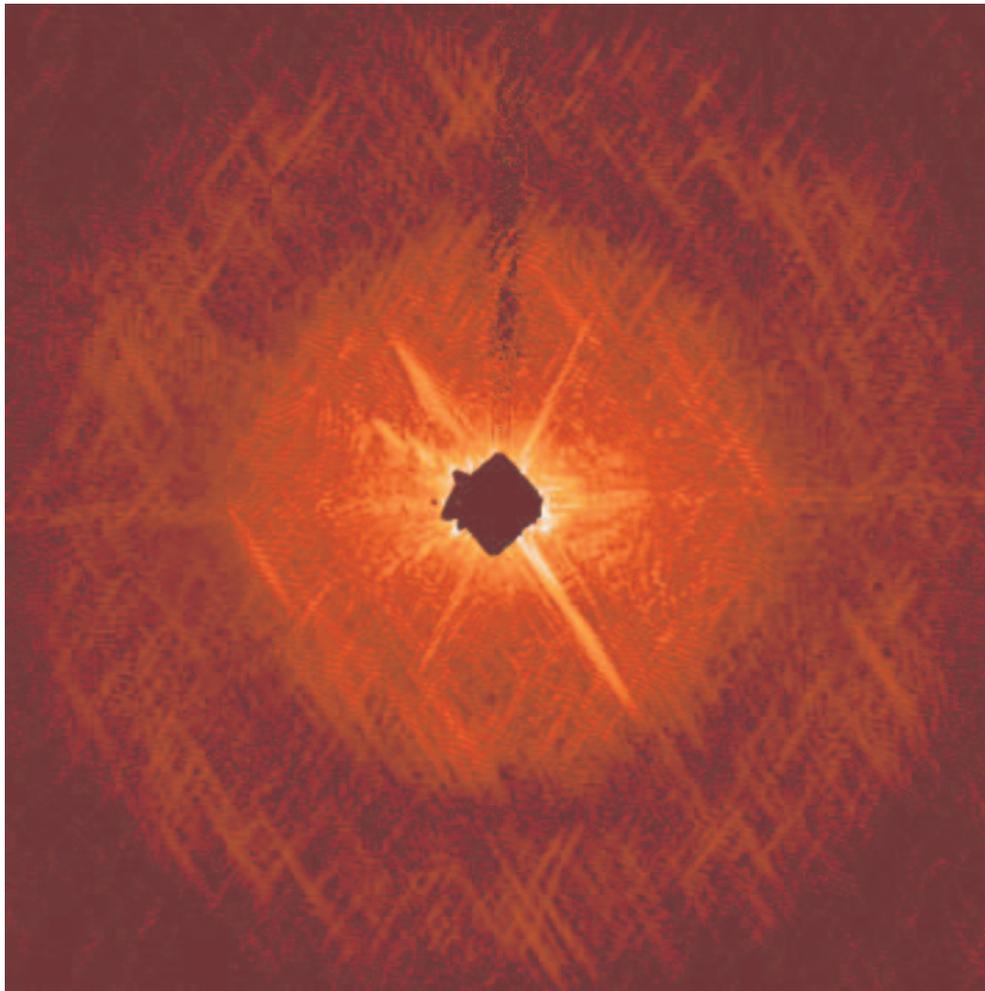


Stony Brook diffraction chamber
(Chris Jacobsen and Janos Kirz)

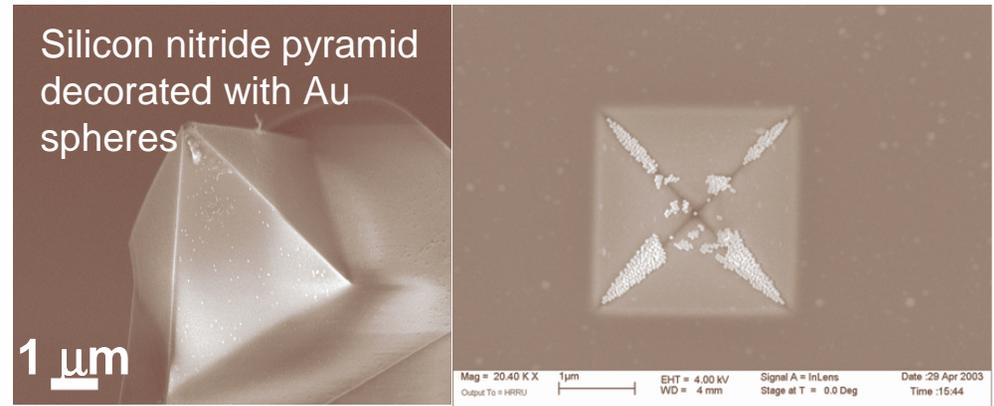


We performed 3D diffraction-imaging experiments

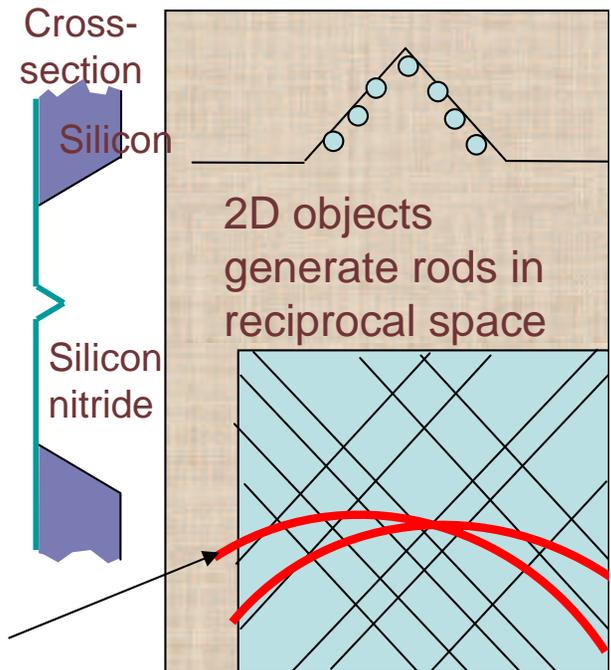
- Complete coverage of reciprocal space by sample rotation
- Use a true 3D object that can be well-characterized by independent means



Diffraction pattern

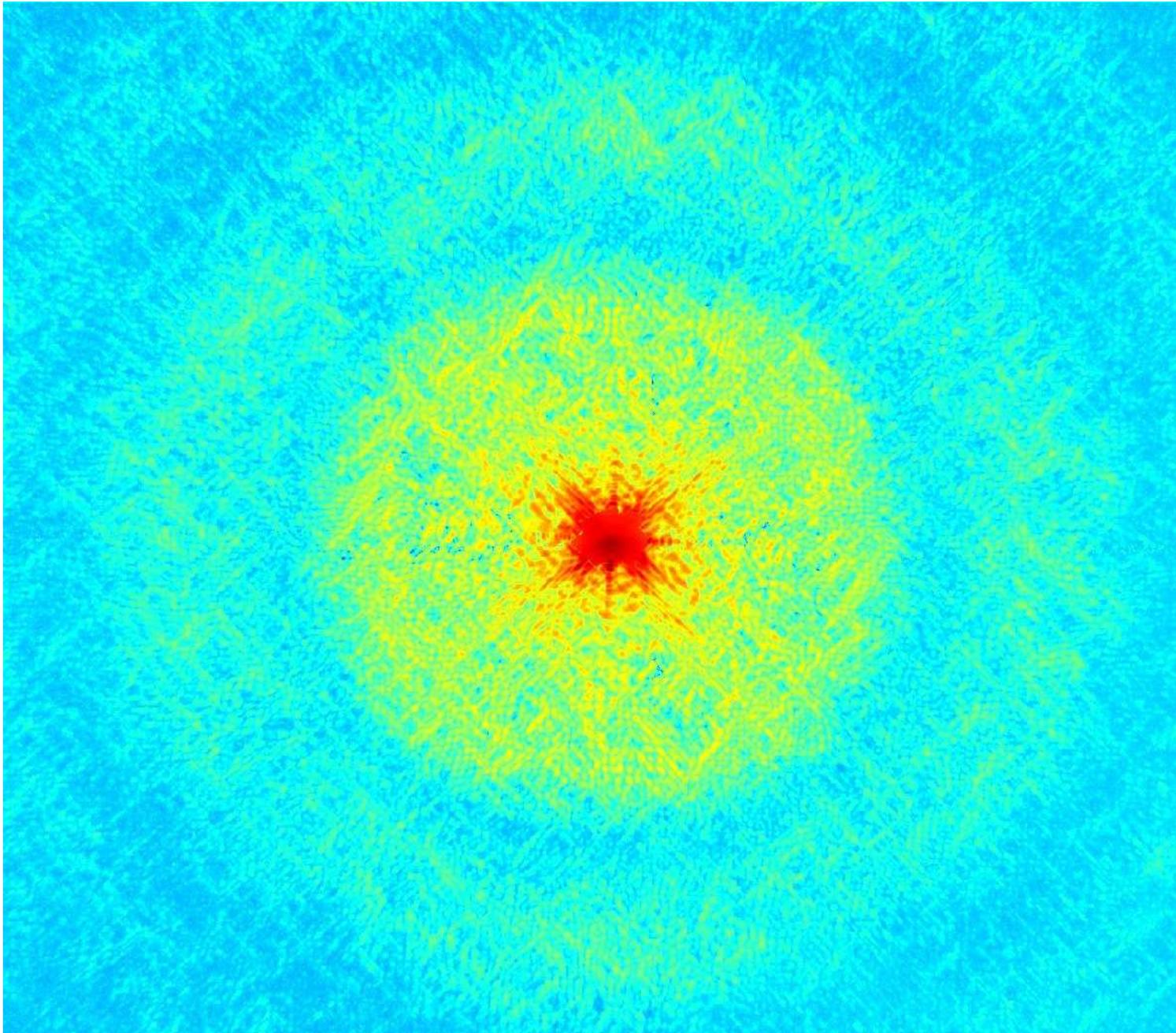


We collected a complete data set with over 140 views with 1° angular spacing.



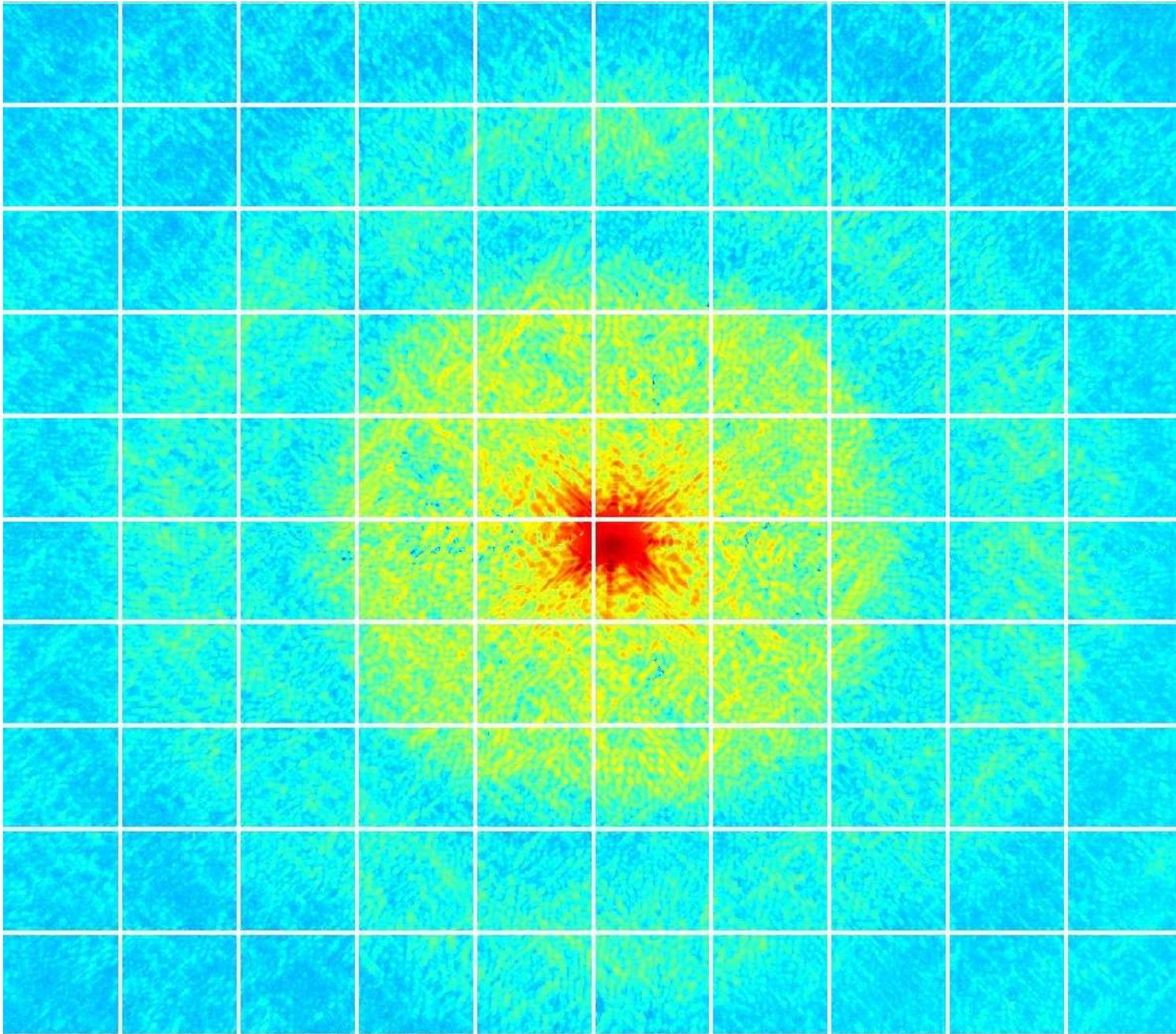
Ewald sphere

Inspection of the data (I)



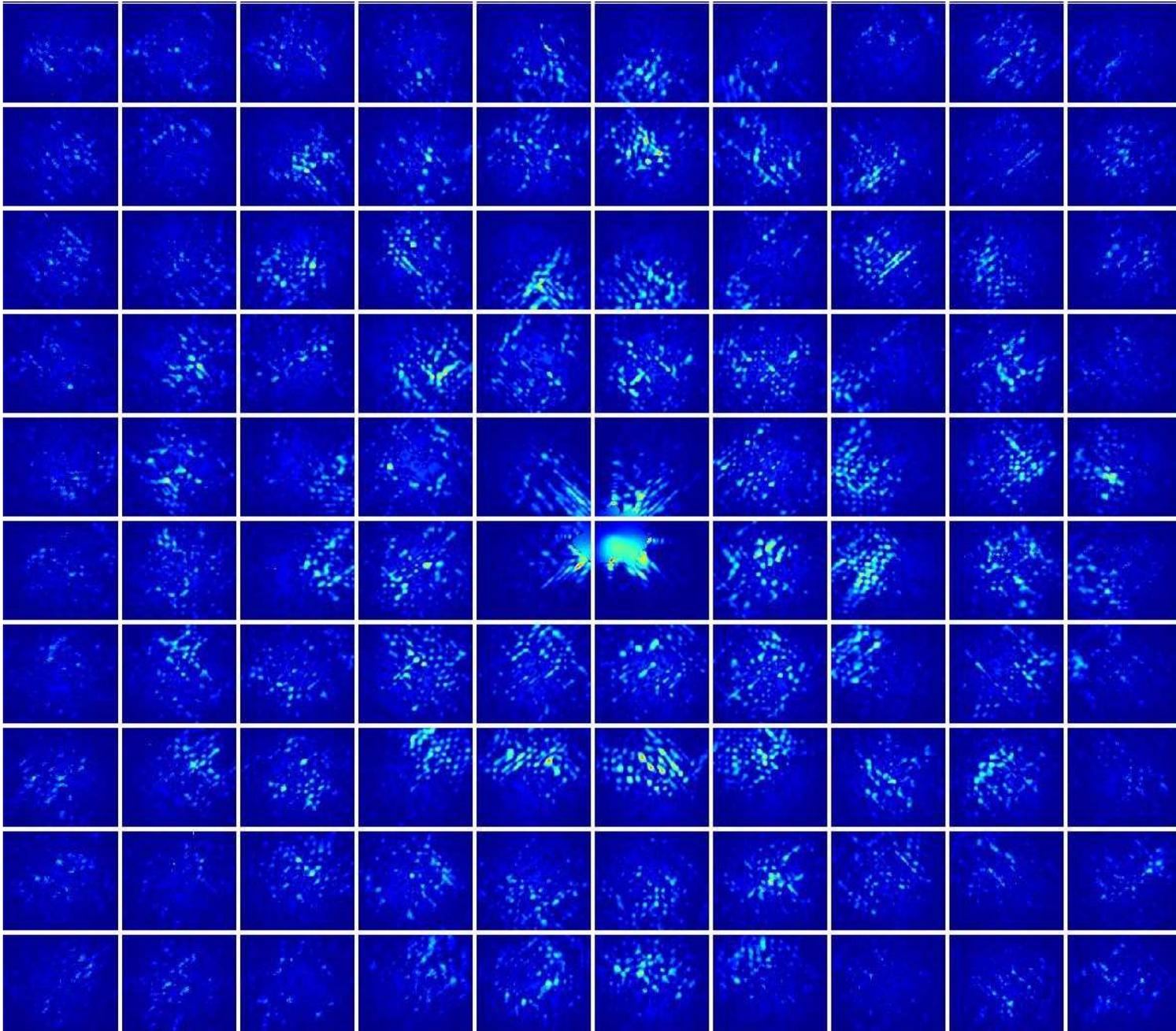
Diffraction
Pattern

Inspection of the data (II)



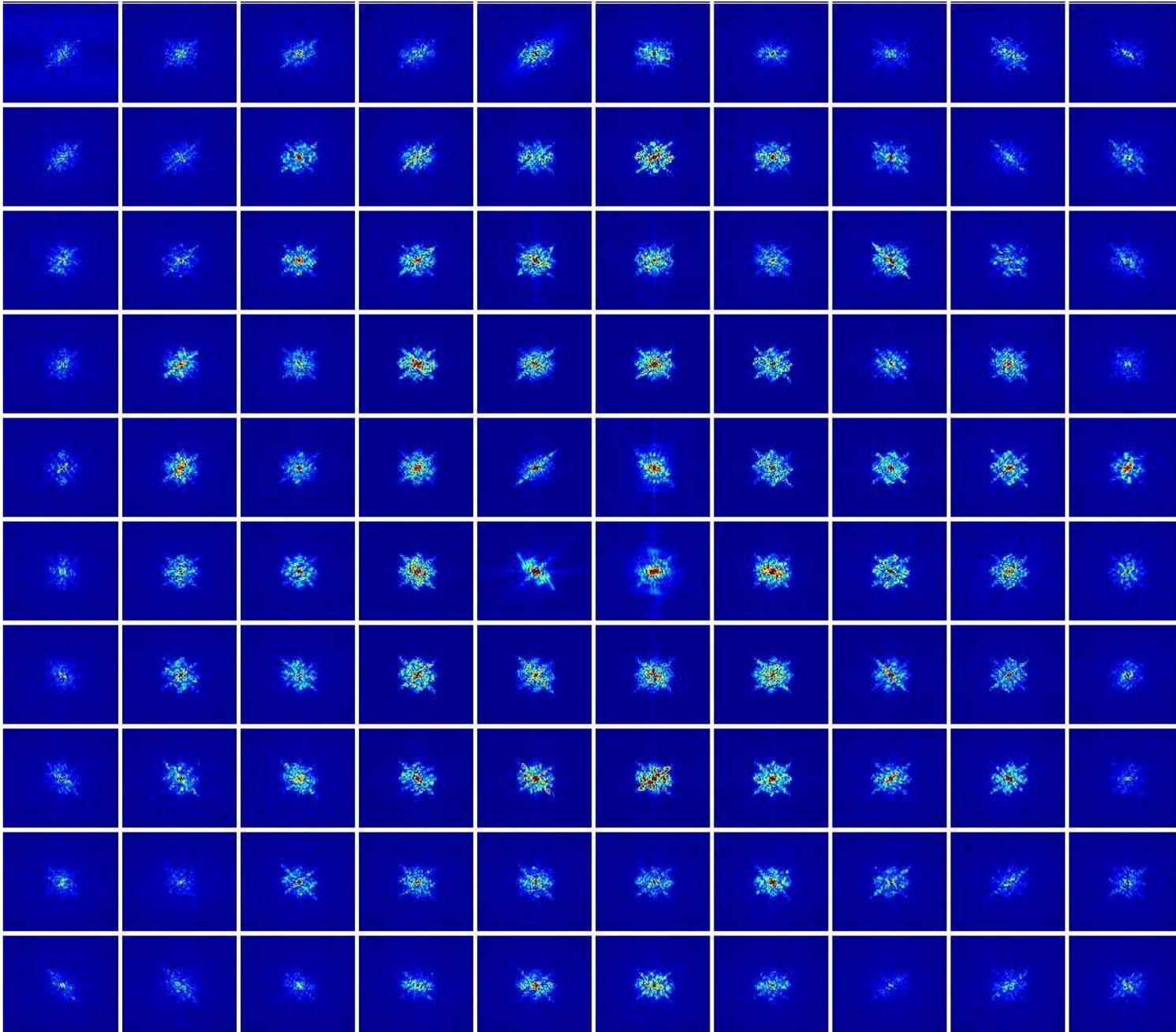
Divide into many pieces

Inspection of the data (III)



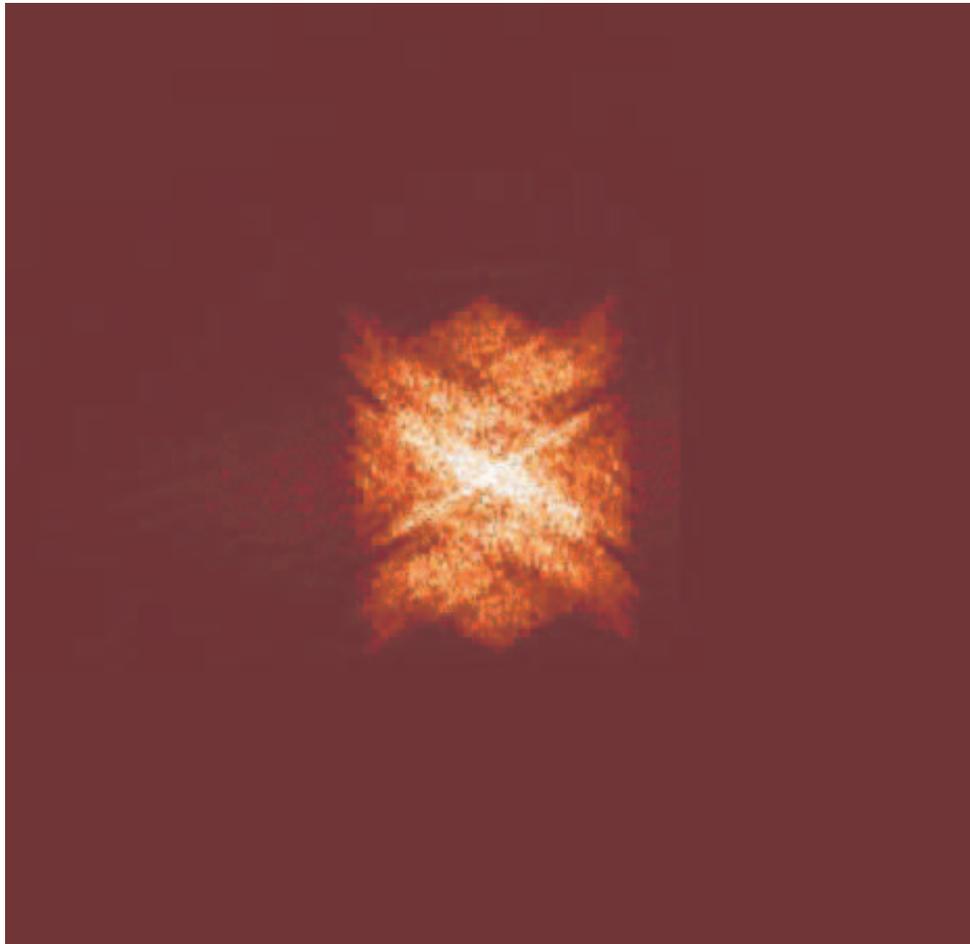
Multiply by
gaussians

Inspection of the data (IV)

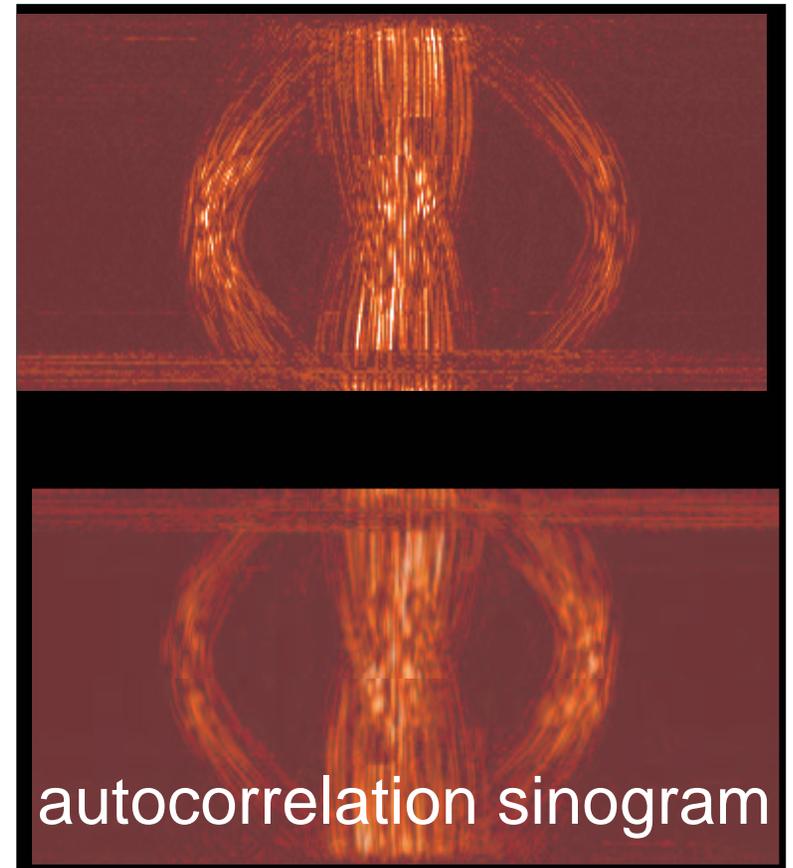


IFFT each
patch

High pass filtered autocorrelation

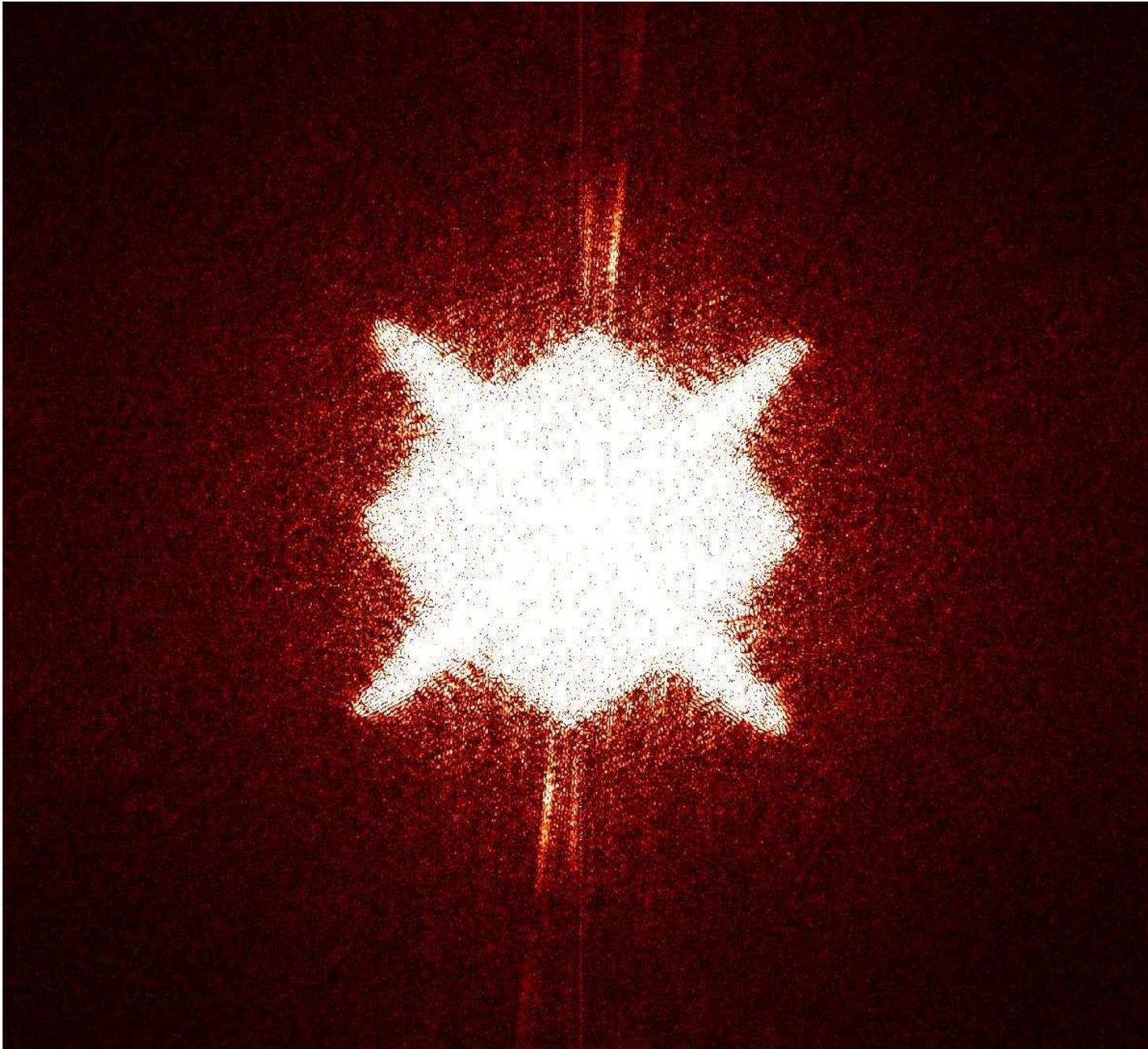


Autocorrelation
(Fourier transform of diffraction intensity)



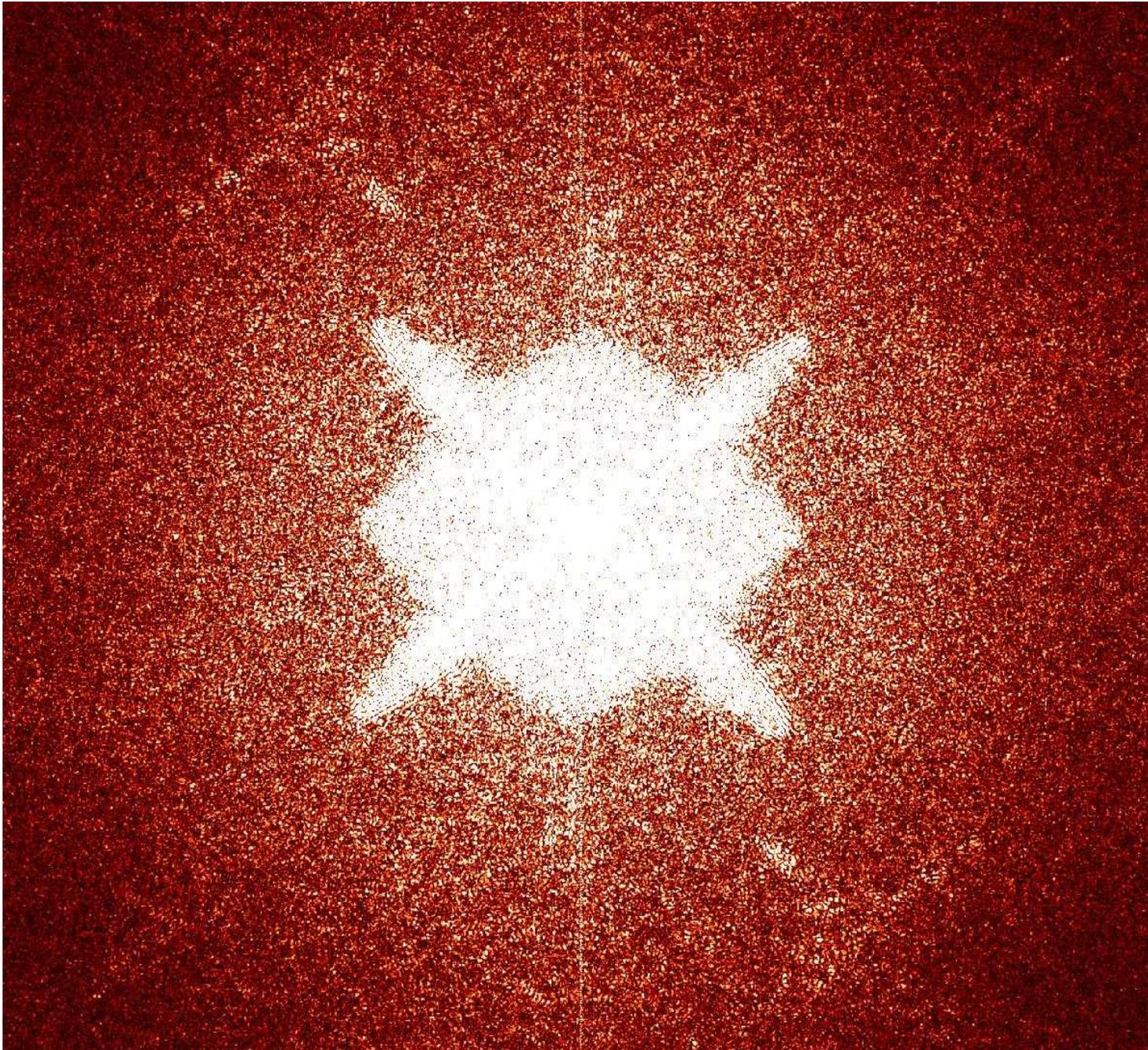
autocorrelation sinogram

Inspection of the data



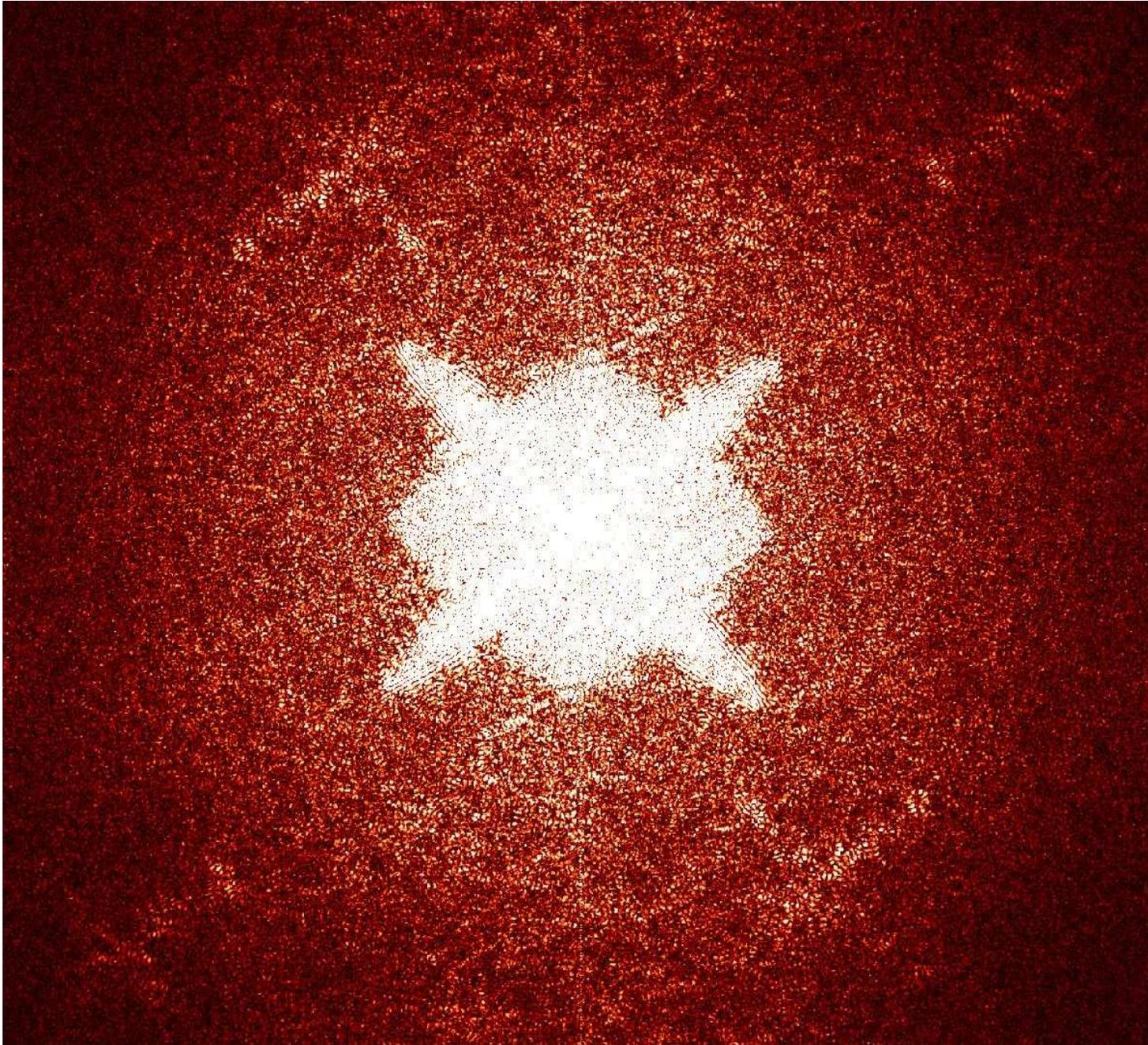
Calculate
Autocorrelation
function

Hologram



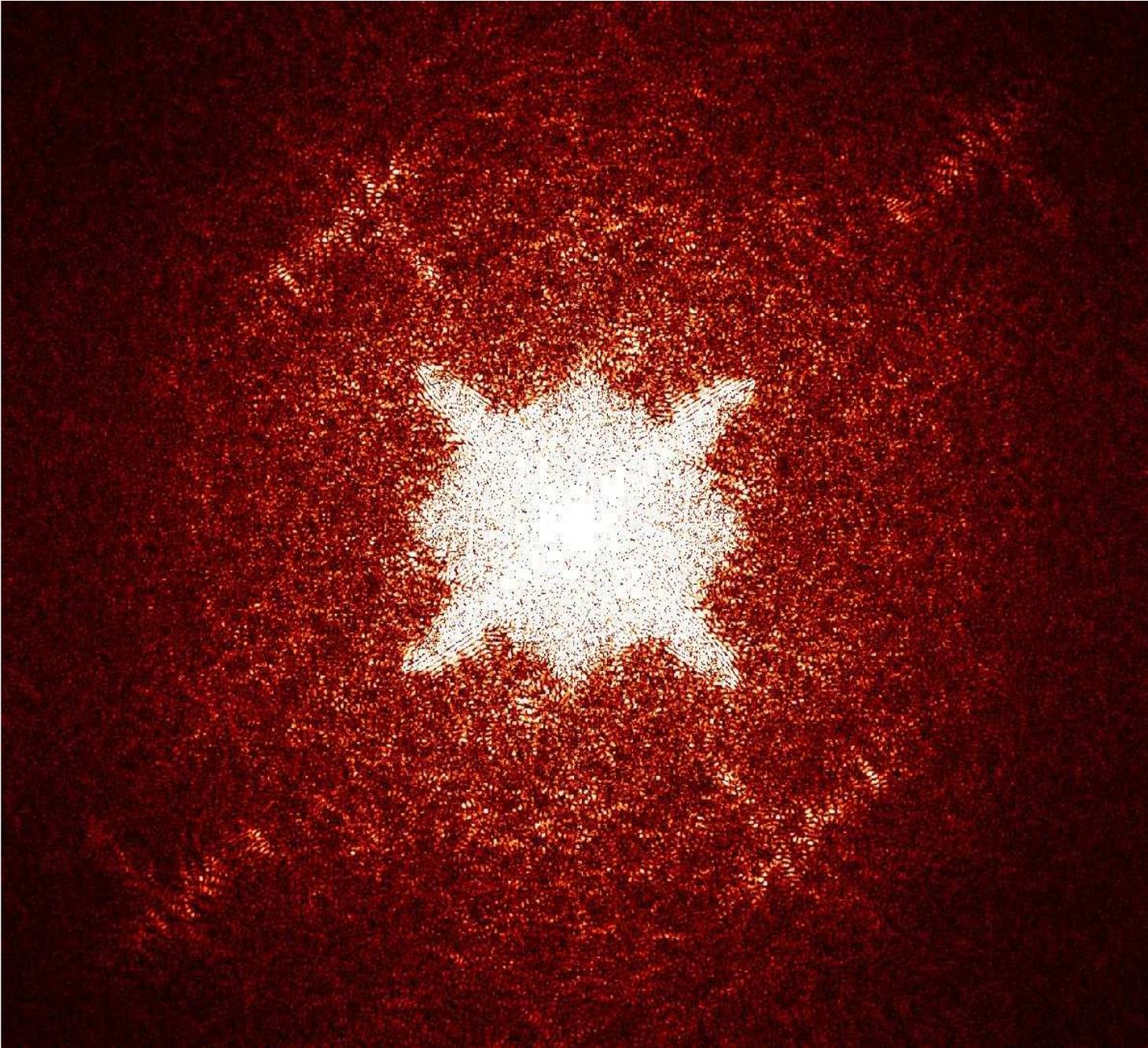
Increase
illumination
area

Hologram



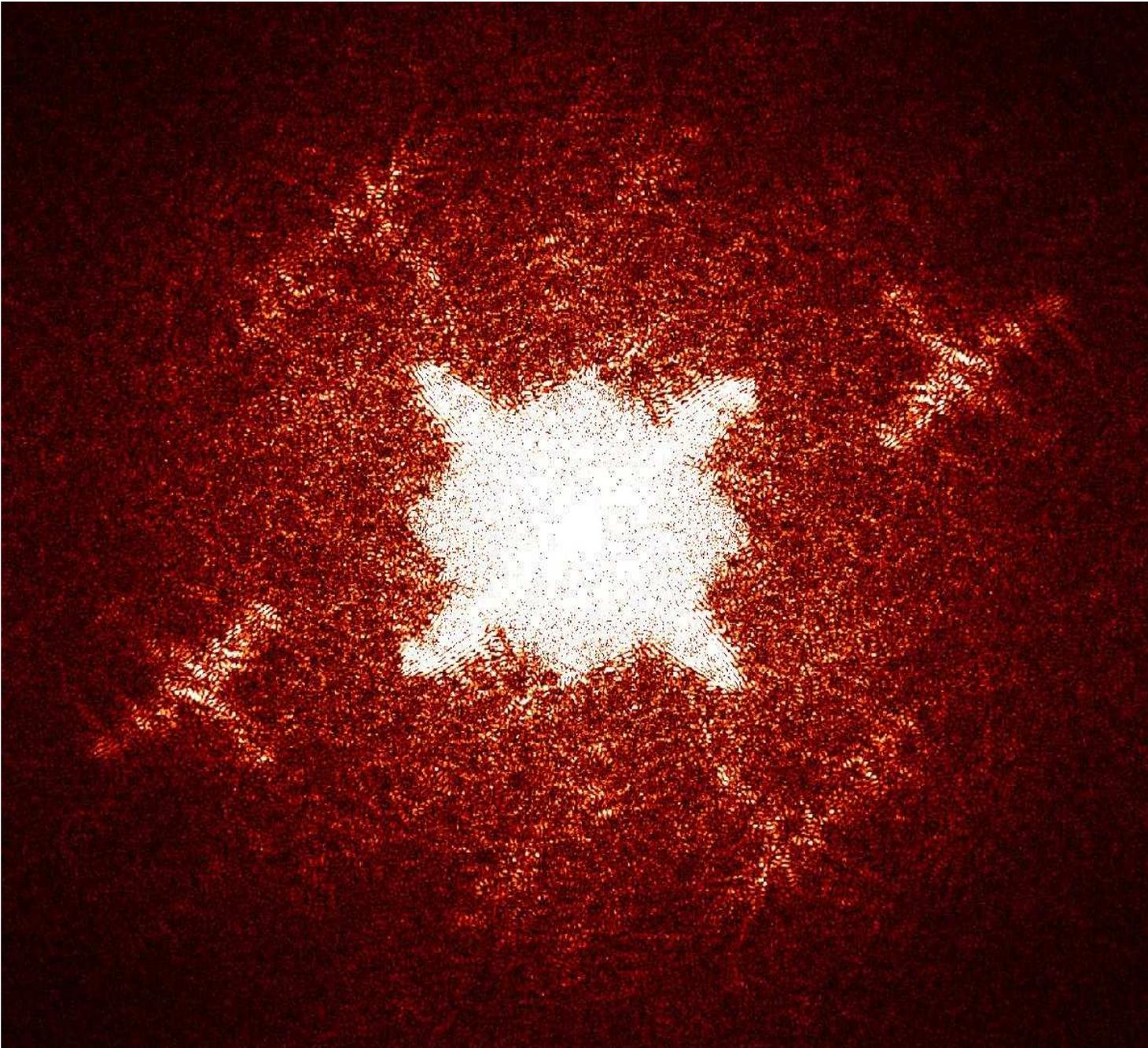
Increase
detector
distance

Hologram



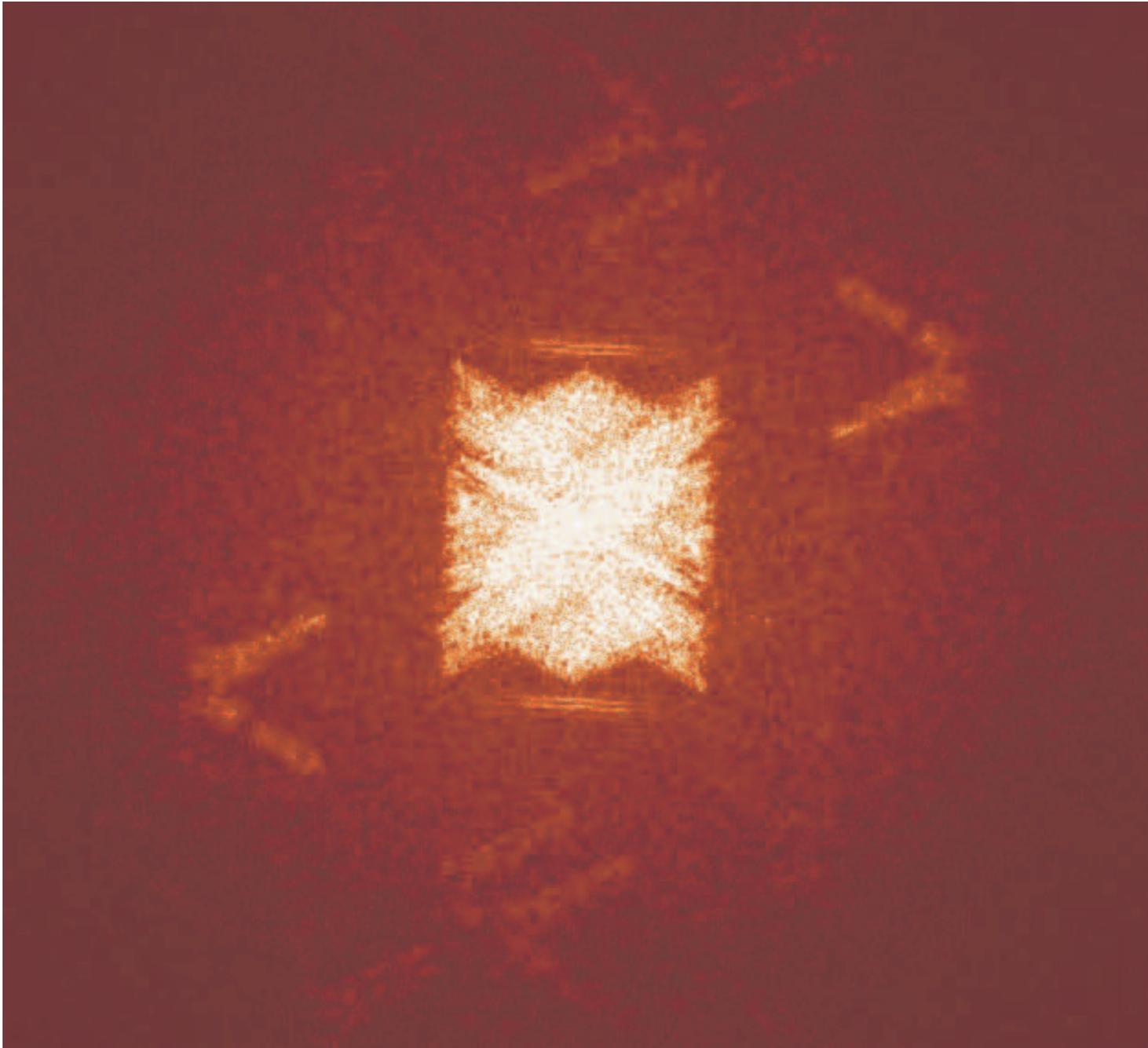
Increase
detector
distance

Hologram



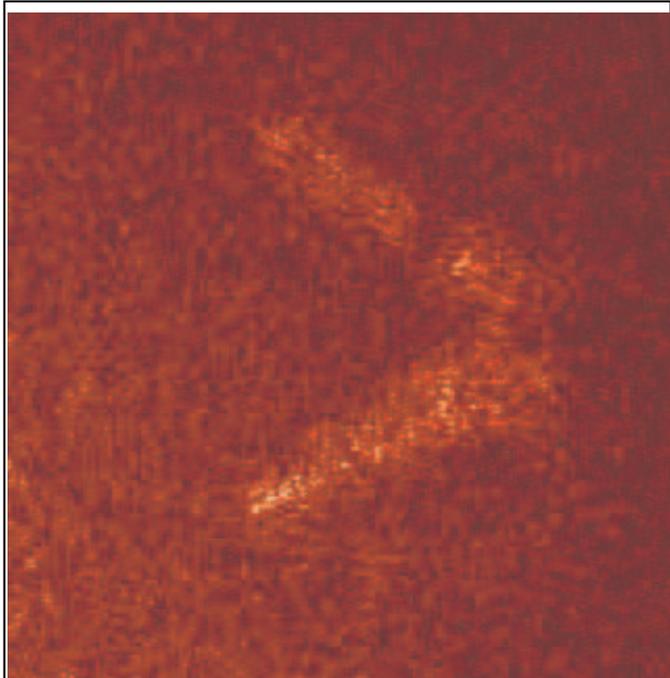
Move beam
center to the
side

Hologram



Rotate

Resolution can be extended by deconvolution



Low-resolution
hologram image

The resolution is limited by the information in the diffraction pattern, not the size of the reference object. We can deconvolve if we know the reference object.

[arXiv:[physics.optics/0405036](https://arxiv.org/abs/physics.optics/0405036)]

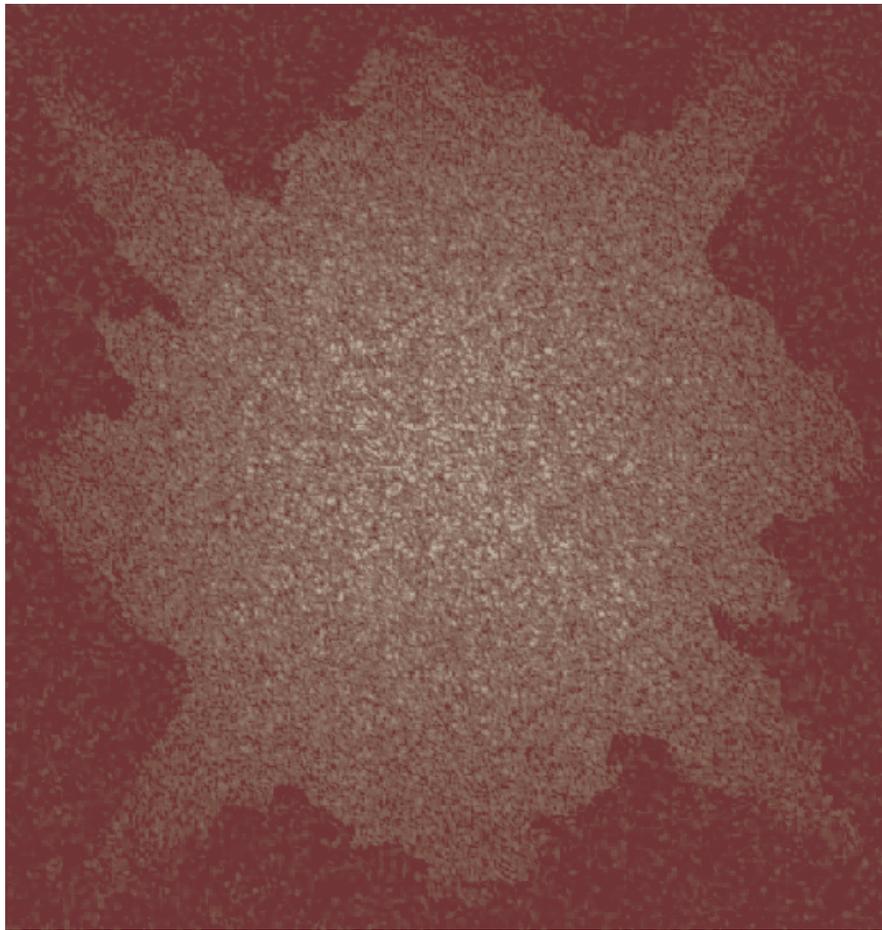


Simulated hologram
image (50 nm
reference object)



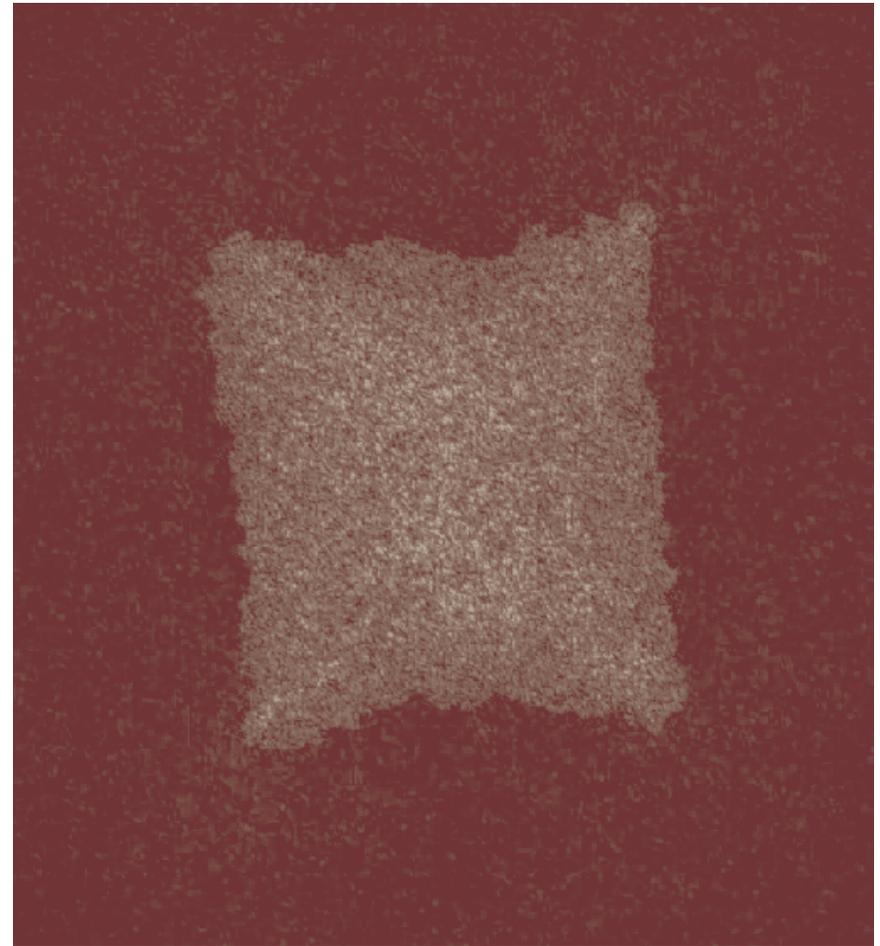
Simulated hologram
image (deconvolved)

Shrinkwrap can reconstruct 2D views of 3D objects



13 degrees rotation

Pyramid
orientation:

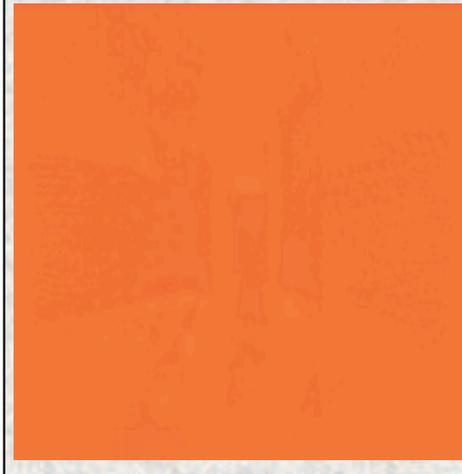


28 degrees rotation

Multiple view reconstructions



Tomography
(from simulation)

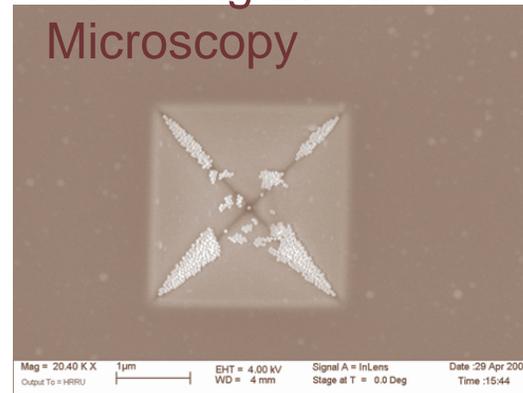


3D-Phase
retrieval

3D- $(1K)^3$ FFT
(109 unknowns)

We need access
to the big LLNL
computers

Scanning Electron
Microscopy



Many error metrics

- real space error (stuff outside support)
- reciprocal space error, (R_{fact})
- difference between reconstructions

In the language of projectors and sets

Reciprocal space error

$$R^n = \frac{\sum_{\mathbf{k}} \left| |\tilde{\rho}_{\text{obs}}(\mathbf{k})| - |\tilde{\rho}^n(\mathbf{k})| \right|^2}{\sum_{\mathbf{k}} |\tilde{\rho}_{\text{obs}}(\mathbf{k})|^2}$$

Real space error

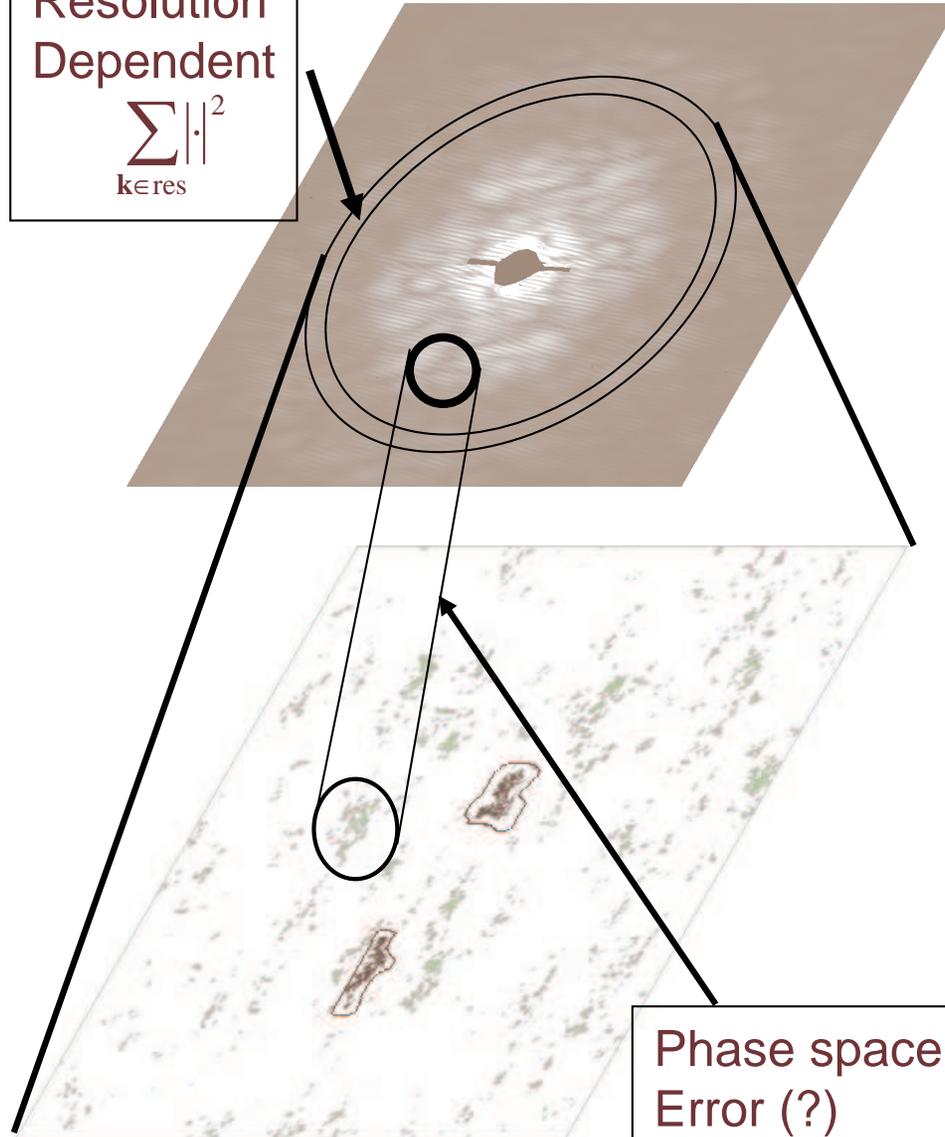
$$\epsilon_s^n = \frac{\sum_{\mathbf{x}} \left| \mathbf{P}_s \rho^n(\mathbf{x}) - \rho^n(\mathbf{x}) \right|^2}{\sum_{\mathbf{x}} \left| \mathbf{P}_s \rho^n(\mathbf{x}) \right|^2}$$

Error between reconstructions

$$\epsilon^{1,2} = \frac{\left\| \rho^1 - \rho^2 \right\|^2}{\left\| \frac{1}{2} (\rho^1 + \rho^2) \right\|^2}$$

Resolution
Dependent

$$\sum_{\mathbf{k} \in \text{res}} |\cdot|^2$$



Phase space
Error (?)

Parseval

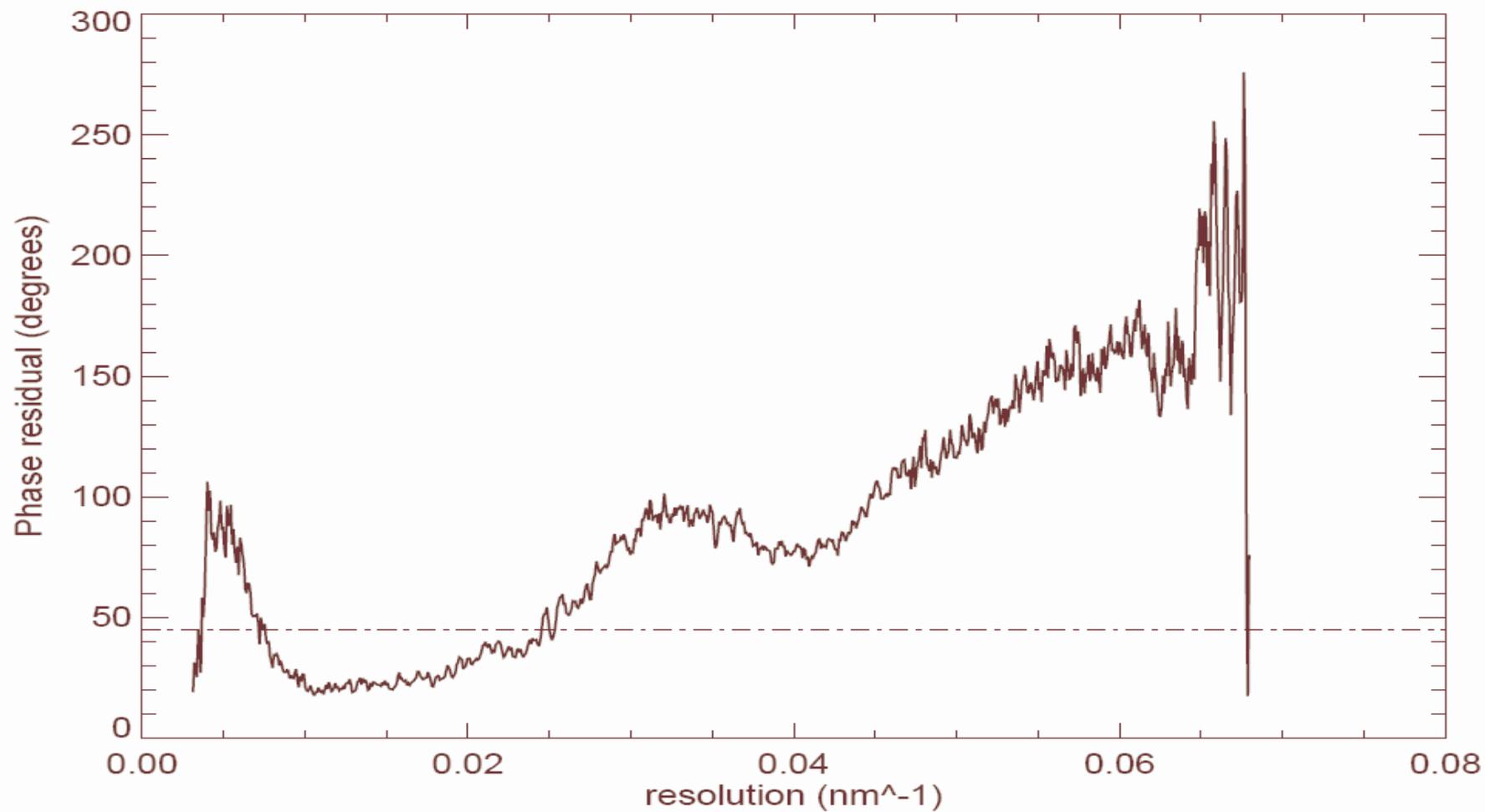
$$\begin{aligned} \|\cdot\|^2 &= \sum_{\mathbf{x}} |\cdot|^2 \\ &= \frac{1}{N} \sum_{\mathbf{k}} |\tilde{\cdot}|^2 \end{aligned}$$

Constraints errors

$$\epsilon_m^n = \frac{\left\| \mathbf{P}_m \rho^n - \rho^n \right\|^2}{\left\| \mathbf{P}_m \rho^n \right\|^2}$$

$$\epsilon_s^n = \frac{\left\| \mathbf{P}_s \rho^n - \rho^n \right\|^2}{\left\| \mathbf{P}_s \rho^n \right\|^2}$$

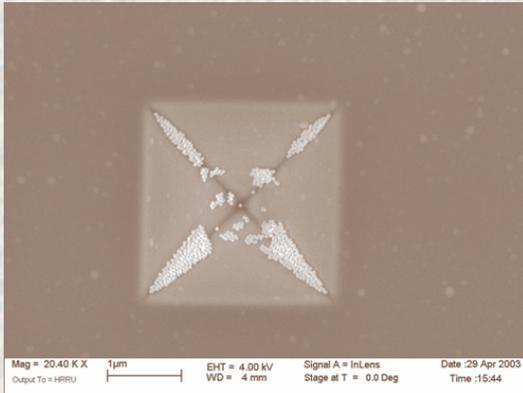
Pyramid data



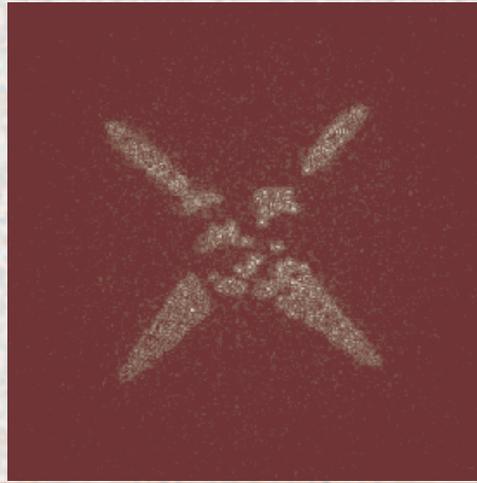
30 seconds exposure

Summary

Scanning Electron Microscopy



2D-Phase retrieval With Shrink-wrap



Tomography (from simulation)



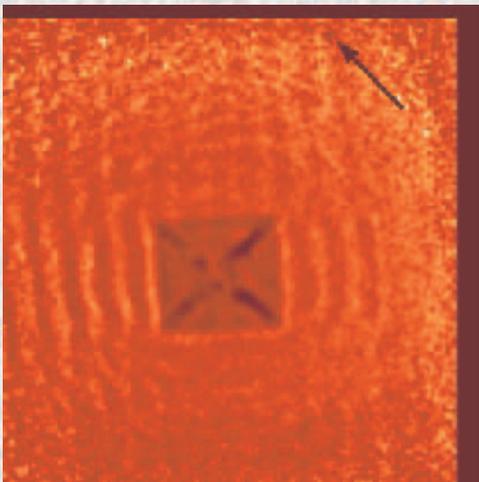
3D-Phase retrieval

3D- $(1K)^3$ FFT
(109 unknowns)

We need access
to the big (LLNL)
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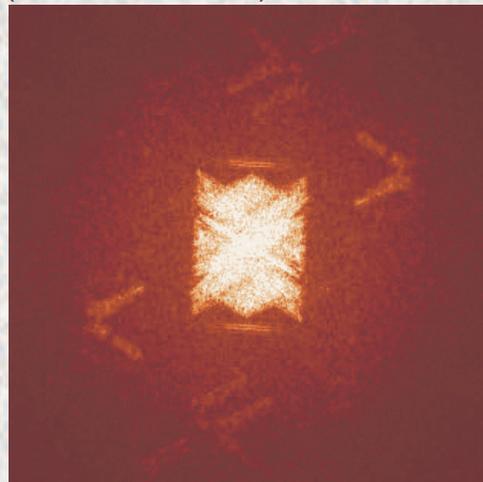
X-ray microscope

The isolated ball acts as a reference for a **Hologram**



Hologram

- Larger illumination area (Weaker intensity)
- Larger camera distance (Lower resolution)



Algorithms being tested:

•Conjugate-gradient based SPEDEN:

Reconstructing single particles from their diffraction patterns, Stefan P. Hau-Riege, Hanna Szöke, Henry N. Chapman, Abraham Szöke, Stefano Marchesini, Alexander Noy, et al.

[arXiv:[physics.optics/0403091](https://arxiv.org/abs/physics/optics/0403091)]

•“Direct Methods”

Experimental lensless soft-X-ray imaging using Direct Methods: phasing diffuse scattering, J. Wu, C. Giacovazzo, S. Marchesini et al.

[arXiv:[physics.optics/0404073](https://arxiv.org/abs/physics/optics/0404073)]

Fourier Transform Holography

Use of extended and prepared reference objects in experimental Fourier transform X-ray holography, H. He, M. Howells, S. Marchesini, H. Chapman, U. Weierstall and J.C.H.Spence.

[arXiv:[physics.optics/0405036](https://arxiv.org/abs/physics/optics/0405036)]